

## Higher Harmonic Anisotropic Flow Measurements of Charged Particles in Pb-Pb Collisions at $\sqrt{s_{NN}} = 2.76$ TeV

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We report on the first measurement of the triangular  $v_3$ , quadrangular  $v_4$ , and pentagonal  $v_5$  charged particle flow in Pb-Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV measured with the ALICE detector at the CERN Large Hadron Collider. We show that the triangular flow can be described in terms of the initial spatial anisotropy and its fluctuations, which provides strong constraints on its origin. In the most central events, where the elliptic flow  $v_2$  and  $v_3$  have similar magnitude, a double peaked structure in the two-particle azimuthal correlations is observed, which is often interpreted as a Mach cone response to fast partons. We show that this structure can be naturally explained from the measured anisotropic flow Fourier coefficients.

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The quark-gluon plasma is a state of matter whose existence at high-energy density is predicted by quantum chromodynamics. The creation of this state of matter in the laboratory and the study of its properties are the main goals of the ultrarelativistic nuclear collision program. One of the experimental observables that is sensitive to the properties of this matter is the azimuthal distribution of particles in the plane perpendicular to the beam direction. When nuclei collide at nonzero impact parameter (noncentral collisions), the geometrical overlap region is anisotropic. This initial spatial asymmetry is converted via multiple collisions into an anisotropic momentum distribution of the produced particles [1].

The azimuthal anisotropy is usually characterized by the Fourier coefficients [2,3]:

$$v_n = \langle \cos[n(\phi - \Psi_n)] \rangle, \quad (1)$$

where  $\phi$  is the azimuthal angle of the particle,  $\Psi_n$  is the angle of the initial state spatial plane of symmetry, and  $n$  is the order of the harmonic. Because the planes of symmetry  $\Psi_n$  are not known experimentally, the anisotropic flow coefficients are estimated from measured correlations between the observed particles. The second Fourier coefficient  $v_2$  is called elliptic flow and has been studied in detail in recent years [4]. Large values of elliptic flow at the LHC were recently observed by the ALICE Collaboration [5].

In a noncentral heavy ion collision the beam axis and the impact parameter define the reaction plane  $\Psi_{RP}$ . Assuming a smooth matter distribution in the colliding nuclei, the plane of symmetry is the reaction plane  $\Psi_n = \Psi_{RP}$  and the

odd Fourier coefficients are zero by symmetry. However, due to fluctuations in the matter distribution, including contributions from fluctuations in the positions of the participating nucleons in the nuclei, the plane of symmetry fluctuates event by event around the reaction plane. This plane of symmetry is determined by the participating nucleons and is therefore called the participant plane  $\Psi_{pp}$  [6]. Event-by-event fluctuations of the spatial asymmetry generate additional odd harmonic symmetry planes  $\Psi_n$ , which are predicted to give rise to the odd harmonics like  $v_3$  and  $v_5$  [7–13].

The large elliptic flow at the Relativistic Heavy Ion Collider (RHIC) [14,15] and at the LHC [5] provides compelling evidence for strongly interacting matter which appears to behave like an almost perfect (inviscid) fluid [16]. Deviations from this ideal case are controlled by the ratio  $\eta/s$  of shear viscosity to entropy density. Because the effect of shear viscosity is to dampen all coefficients, with a larger decrease for higher order coefficients [12,17], it has been argued that the magnitude and transverse momentum dependence of the coefficients  $v_3$  and  $v_5$  is a more sensitive measure of  $\eta/s$  [11]. Therefore a measurement of these Fourier coefficients at the LHC provides strong constraints on the initial geometry, its fluctuations, as well as on the shear viscosity to entropy density ratio.

In this Letter we report the first measurement of the anisotropic flow coefficients  $v_3$ ,  $v_4$ , and  $v_5$  of charged particles in Pb-Pb collisions at the center of mass energy per nucleon pair  $\sqrt{s_{NN}} = 2.76$  TeV, with the ALICE detector [18–20]. The data were recorded in November 2010 during the first run with heavy ions at the LHC.

For this analysis the ALICE inner tracking system (ITS) and the time projection chamber (TPC) were used to reconstruct charged particle tracks. The VZERO counters and the silicon pixel detector (SPD) were used for the trigger. The VZERO counters are two scintillator arrays providing both amplitude and timing information, covering

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the pseudorapidity range  $2.8 < \eta < 5.1$  (VZERO-A) and  $-3.7 < \eta < -1.7$  (VZERO-C). The SPD is the innermost part of the ITS, consisting of two cylindrical layers of hybrid silicon pixel assemblies covering the range of  $|\eta| < 2.0$  and  $|\eta| < 1.4$  for the inner and outer layer, respectively. The minimum-bias interaction trigger required the following three conditions [21]: (i) two pixel chip hits in the outer layer of the silicon pixel detectors, (ii) a signal in VZERO-A, and (iii) a signal in VZERO-C. Deflection of neutral recoils, which is sensitive to the directed flow of spectators, is measured with two neutron zero degree calorimeters (ZDCs) installed on each side, 114 m from the interaction point. Only events with a vertex found within 7 cm from the center of the detector along the beam line were used in the analysis. This is to ensure a uniform acceptance in the central pseudorapidity region  $|\eta| < 0.8$ . An event sample of  $5 \times 10^6$  Pb-Pb collisions passed the selection criteria and was analyzed as a function of collision centrality, determined by cuts on the VZERO multiplicity as described in [21]. Based on the strong correlation between the collision centrality determined by the ZDC, TPC, SPD, and VZERO detectors, the resolution in centrality is found to be  $< 0.5\%$  rms for the most central collisions (0%–5%), increasing towards 2% rms for peripheral collisions (e.g., 70%–80%). This resolution is also in agreement with our Monte Carlo (MC) Glauber [22] studies.

The analysis was, as in [5], performed using tracks measured with only the TPC and for tracks using the ITS and TPC. These two measurements have very different acceptance and efficiency corrections, and provide an estimate of a possible residual bias which may be present, after correction, for small values of the harmonics [23]. For both measurements, charged particles were selected with high reconstruction efficiency and minimal contamination from photon conversions and secondary charged particles produced in the detector material as described in [5]. From Monte Carlo simulations of HIJING [24] events using a GEANT3 [25] detector simulation and event reconstruction, the estimated contamination is less than 6% at  $p_t = 0.2$  GeV/ $c$  and drops below 1% at  $p_t > 1$  GeV/ $c$ . In this Letter we present the results obtained using the TPC stand-alone tracks, because of the smaller corrections for the azimuthal acceptance.

We report the anisotropic flow coefficients  $v_n$  obtained from two-particle correlations and from a four-particle cumulant method [26], denoted  $v_n\{2\}$  and  $v_n\{4\}$ , respectively. To calculate the four-particle cumulants we used the method proposed in [23]. The  $v_n\{2\}$  and  $v_n\{4\}$  measurements have different sensitivity to flow fluctuations and contributions from nonflow. The nonflow contribution arises from correlations between the particles unrelated to the initial geometry. The contribution from flow fluctuations is positive for  $v_n\{2\}$  while it is negative for  $v_n\{4\}$  [27]. Because the odd harmonics are expected to be completely

due to event-by-event fluctuations in the initial spatial geometry, the comparison of these two- and four-particle cumulants provides a strong constraint on the initial spatial geometry fluctuations.

The nonflow contribution to the two-particle correlations is not known and might be significant. We utilize four methods to study and correct for nonflow contributions to the  $v_n\{2\}$  coefficients. First, we compare  $v_n\{2\}$  for like and unlike charge-sign combinations since they have different contributions from resonance decay and jet fragmentation. Second, we used different pseudorapidity gap requirements between the two particles since larger gaps reduce the nonflow contributions. Third, we utilize HIJING (a perturbative quantum chromodynamics inspired model which does not include flow) to estimate these contributions, and, finally, we estimate the nonflow from the correlations measured in proton-proton collisions. All of these methods indicate that nonflow effects are smaller than 10%. In this Letter we use the dependence of the correlations on pseudorapidity distance between particles as an estimate of nonflow.

Figure 1(a) shows  $v_2$ ,  $v_3$ , and  $v_4$  integrated over the  $p_t$  range  $0.2 < p_t < 5.0$  GeV/ $c$  as a function of centrality. The  $v_2\{2\}$ ,  $v_3\{2\}$ , and  $v_4\{2\}$  are shown for particles with  $|\Delta\eta| > 1.0$  and corrected for the estimated remaining nonflow contribution based on the correlation measured in HIJING. The total systematic uncertainty is shown as a band and fully includes this residual correction. The measured  $v_3$  is smaller than  $v_2$  and does not depend strongly on centrality. The  $v_3$  is compatible with predictions for Pb-Pb collisions from a hydrodynamic model calculation with Glauber initial conditions and  $\eta/s = 0.08$  and larger than for MC-KLN (Kharzeev-Levin-Nardi) color glass condensate (CGC) [28] initial conditions with  $\eta/s = 0.16$  [11], suggesting a small value of  $\eta/s$  for the matter created in these collisions. The  $v_3\{4\}$  is about a factor 2 smaller than the two-particle measurement which can, as explained in [29], be understood if  $v_3$  originates predominantly from event-by-event fluctuations of the initial spatial geometry. For these event-by-event fluctuations of the spatial geometry, the symmetry plane  $\Psi_3$  is expected to be uncorrelated (or correlated very weakly [30]) with the reaction plane  $\Psi_{RP}$  and with  $\Psi_2$ . We evaluate the correlations between  $\Psi_3$  and  $\Psi_{RP}$  using the first-order event plane from the ZDC via  $v_{3/\Psi_{RP}} = \langle \cos(3\phi_1 - 3\Psi_{RP}) \rangle$  and the correlation between  $\Psi_3$  and  $\Psi_2$  with a five-particle correlator  $\langle \cos(3\phi_1 + 3\phi_2 - 2\phi_3 - 2\phi_4 - 2\phi_5) \rangle / v_2^3 = v_{3/\Psi_2}^2$ . In Fig. 1(a)  $v_{3/\Psi_{RP}}$  and  $v_{3/\Psi_2}^2$  are shown as a function of centrality. These correlations are indeed, within uncertainties, consistent with zero, as expected from a triangular flow that originates predominantly from event-by-event fluctuations of the initial spatial geometry.

To investigate the role of viscosity further we calculate the ratios  $v_2/\varepsilon_2$  and  $v_3/\varepsilon_3$ , where  $\varepsilon_2$  and  $\varepsilon_3$  are the ellipticity and triangularity of the initial spatial geometry, defined by

$$\varepsilon_n = -\frac{\langle r^2 \cos n(\phi - \Psi_n) \rangle}{\langle r^2 \rangle}, \quad (2)$$

where the brackets denote an average which traditionally is taken over the position of participating (wounded) nucleons in a Glauber model [22].

Under the assumption that  $v_n$  is proportional to  $\varepsilon_n$ ,  $v_n\{2\}$  is proportional to  $\varepsilon_n\{2\}$  [27]. Figure 1(b) shows the ratios  $v_n/\varepsilon_n$  for eccentricities calculated with a Glauber and a MC-KLN CGC [28] model, denoted by  $\varepsilon_n^W\{2\}$  and  $\varepsilon_n^{\text{CGC}}\{2\}$ , respectively. We find that for a Glauber model the magnitude of  $v_3\{2\}/\varepsilon_3\{2\}$  is smaller than  $v_2\{2\}/\varepsilon_2\{2\}$ , which would indicate significant viscous corrections. For

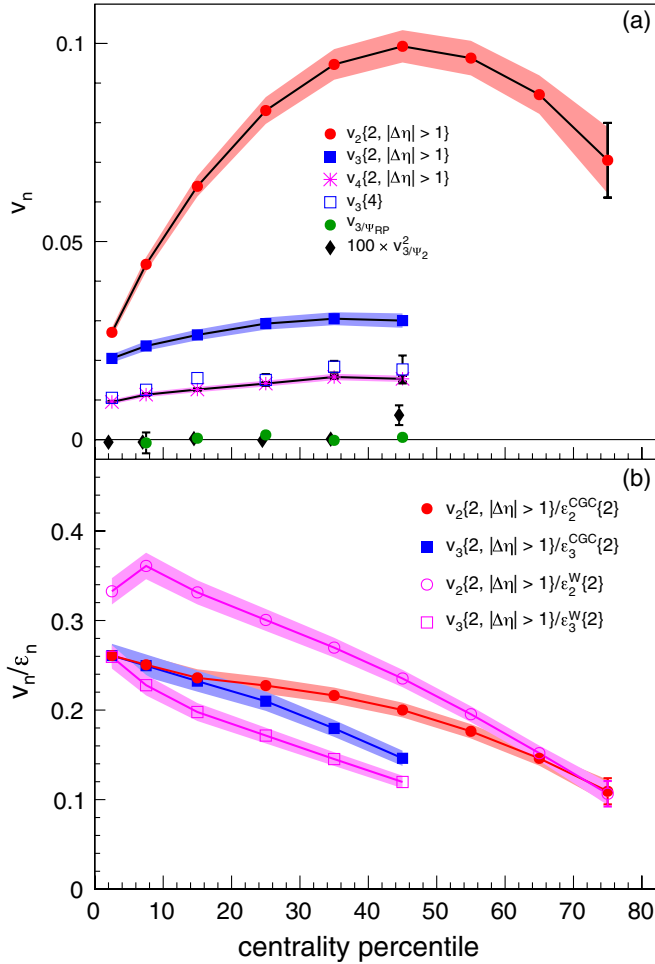


FIG. 1 (color online). (a)  $v_2$ ,  $v_3$ , and  $v_4$  integrated over the  $p_t$  range  $0.2 < p_t < 5.0$  GeV/c as a function of event centrality, with the more central (peripheral) collisions shown on the left-(right-)hand side, respectively. Full and open squares show  $v_3\{2\}$  and  $v_3\{4\}$ , respectively. In addition we show  $v_{3/\Psi_2}^2$  and  $v_{3/\Psi_{\text{RP}}}$ , which represent the triangular flow measured relative to the second order event plane and the reaction plane, respectively (for the definitions, see text). (b)  $v_2\{2, |\Delta\eta| > 1\}$  and  $v_3\{2, |\Delta\eta| > 1\}$  divided by the corresponding eccentricity versus centrality percentile for Glauber [22] and MC-KLN CGC [28] initial conditions.

MC-KLN CGC calculations the ratios  $v_2\{2\}/\varepsilon_2\{2\}$  and  $v_3\{2\}/\varepsilon_3\{2\}$  are almost equal for the most central collisions, as expected for an almost ideal fluid [11]. In addition, we notice that the ratio  $v_3\{2\}/\varepsilon_3\{2\}$  decreases faster than  $v_2\{2\}/\varepsilon_2\{2\}$  toward more peripheral collisions, which is expected due to larger viscous corrections to  $v_3$ .

The centrality dependence of the triangular flow differs significantly from that of elliptic flow. This might be due to two reasons: either the centrality dependence of the spatial ellipticity and triangularity are different and/or the viscous effects are different. However, in a small centrality range, such as 0%–5%, viscous effects do not change much and there one might be directly sensitive to the change in the initial spatial geometry. Our calculations show that even in this small centrality range, the ratio  $\varepsilon_2/\varepsilon_3$  changes significantly, which allows us to investigate further the geometrical origin of elliptical and triangular flow. In Fig. 2  $v_2\{2\}$  and  $v_3\{2\}$  are plotted in 1% centrality bins for the 5% most central collisions. We observe that  $v_3\{2\}$  does not change much versus centrality (as would be expected if  $v_3$  is dominated by event-by-event fluctuations of the initial geometry) while  $v_2\{2\}$  increases by about 60%. We compare this dependence of  $v_n\{2\}$  to the centrality dependence of the eccentricities  $\varepsilon_n\{2\}$  for initial conditions from MC-KLN CGC and Monte Carlo Glauber model. We observe that the weak dependence of  $v_3\{2\}$  is described by both calculations while the relative strong dependence of  $v_2\{2\}$  on centrality is only described for the MC-KLN CGC initial conditions.

The harmonics  $v_2\{2\}$ ,  $v_3\{2\}$ ,  $v_4\{2\}$ , and  $v_5\{2\}$  as a function of transverse momentum are shown for the 30%–40%, 0%–5%, and 0%–2% centrality classes in Fig. 3. For the 30%–40% centrality class the results are compared to hydrodynamic predictions using Glauber initial conditions for different values of  $\eta/s$  [31]. We observe that, at low  $p_t$ , the different  $p_t$  dependence of  $v_2$  and  $v_3$  is described well by these hydrodynamic predictions. However, the

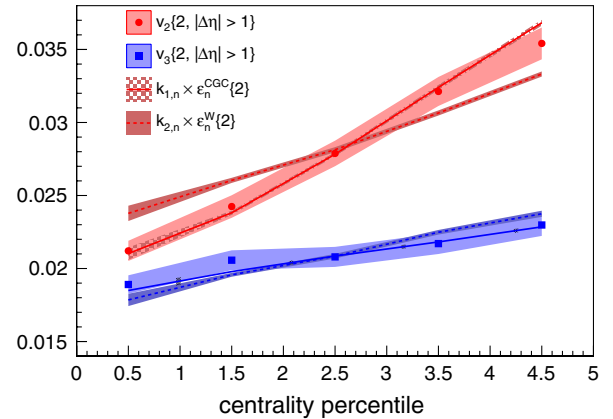


FIG. 2 (color online).  $v_2$  and  $v_3$  as a function of centrality for the 5% most central collisions compared to calculations of the spatial eccentricities,  $\varepsilon_n^W\{2\}$  and  $\varepsilon_n^{\text{CGC}}\{2\}$ . The eccentricities have been scaled to match the 2%–3% data using  $k_1$  and  $k_2$ .

magnitude of  $v_2(p_t)$  is better described by  $\eta/s = 0$  while for  $v_3(p_t)$   $\eta/s = 0.08$  provides a better description. We anticipate future comparisons utilizing MC-KLN initial conditions.

For central collisions 0%–5% we observe that at  $p_t \approx 2$  GeV/c  $v_3$  becomes equal to  $v_2$  and at  $p_t \approx 3$  GeV/c  $v_4$  also reaches the same magnitude as  $v_2$  and  $v_3$ . For more central collisions 0%–2%, we observe that  $v_3$  becomes equal to  $v_2$  at lower  $p_t$  and reaches significantly larger

values than  $v_2$  at higher  $p_t$ . The same is true for  $v_4$  compared to  $v_2$ .

We compare the structures found with azimuthal correlations between triggered and associated particles to those described by the measured  $v_n$  components. The two-particle azimuthal correlations are measured by calculating

$$C(\Delta\phi) \equiv \frac{N_{\text{mixed}}}{N_{\text{same}}} \frac{dN_{\text{same}}/d\Delta\phi}{dN_{\text{mixed}}/d\Delta\phi}, \quad (3)$$

where  $\Delta\phi = \phi_{\text{trig}} - \phi_{\text{assoc}}$ .  $dN_{\text{same}}/d\Delta\phi$  ( $dN_{\text{mixed}}/d\Delta\phi$ ) is the number of associated particles as function of  $\Delta\phi$  within the same (different) event, and  $N_{\text{same}}$  ( $N_{\text{mixed}}$ ) the total number of associated particles in  $dN_{\text{same}}/d\Delta\phi$  ( $dN_{\text{mixed}}/d\Delta\phi$ ). Figure 4 shows the azimuthal correlation observed in very central collisions 0%–1%, for trigger particles in the range  $2 < p_t < 3$  GeV/c with associated particles in  $1 < p_t < 2$  GeV/c for pairs in  $|\Delta\eta| > 1$ . We observe a clear doubly peaked correlation structure centered opposite to the trigger particle. This feature has been observed at lower energies in broader centrality bins [32,33], but only after subtraction of the elliptic flow component. This two-peak structure has been interpreted as an indication for various jet-medium modifications (i.e., Mach cones) [32,33] and more recently as a manifestation of triangular flow [10–13]. We therefore compare the azimuthal correlation shape expected from  $v_2$ ,  $v_3$ ,  $v_4$ , and  $v_5$  evaluated at corresponding transverse momenta with the measured two-particle azimuthal triggered correlation and find that the combination of these harmonics gives a natural description of the observed correlation structure on the away side.

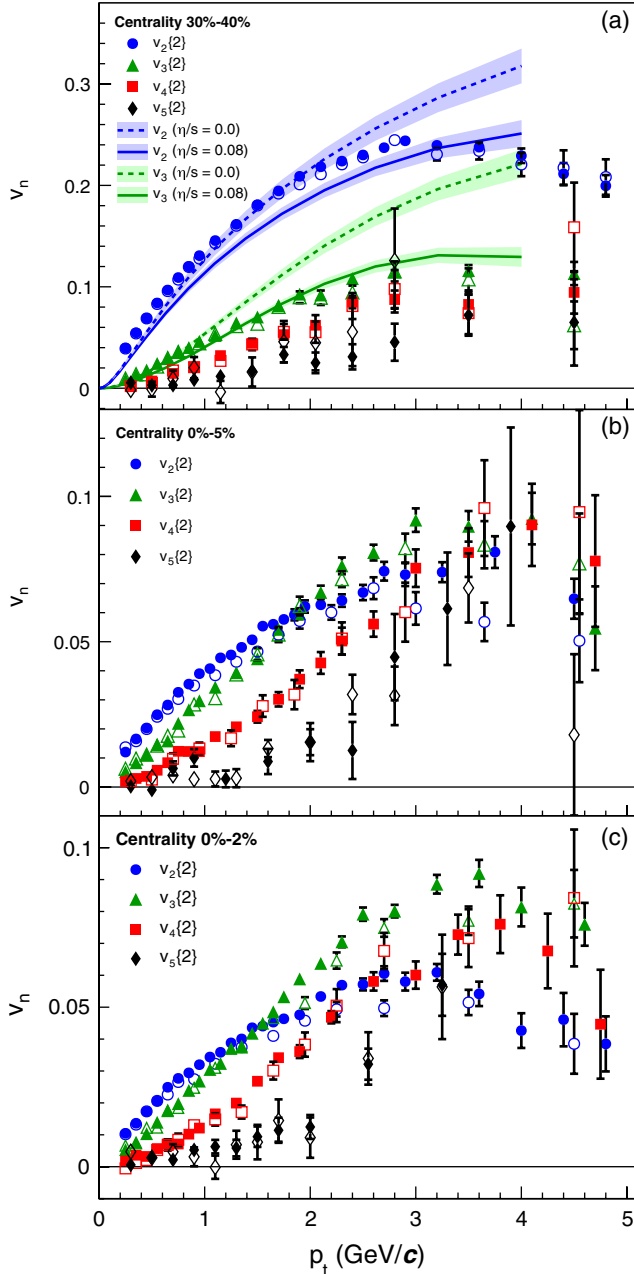


FIG. 3 (color online).  $v_2$ ,  $v_3$ ,  $v_4$ ,  $v_5$  as a function of transverse momentum and for three event centralities. The full and open symbols are for  $\Delta\eta > 0.2$  and  $\Delta\eta > 1.0$ , respectively. (a) 30%–40% compared to hydrodynamic model calculations, (b) 0%–5% centrality percentile, (c) 0%–2% centrality percentile.

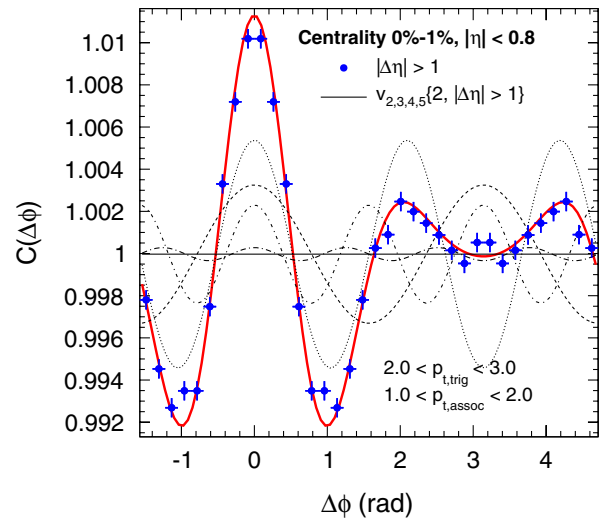


FIG. 4 (color online). The two-particle azimuthal correlation, measured in  $0 < \Delta\phi < \pi$  and shown symmetrized over  $2\pi$ , between a trigger particle with  $2 < p_t < 3$  GeV/c and an associated particle with  $1 < p_t < 2$  GeV/c for the 0%–1% centrality class. The solid red line shows the sum of the measured anisotropic flow Fourier coefficients  $v_2$ ,  $v_3$ ,  $v_4$ , and  $v_5$  (dashed lines).



In summary, we have presented the first measurement at the LHC of triangular  $v_3$ , quadrangular  $v_4$ , and pentagonal particle flow  $v_5$ . We have shown that the triangular flow and its fluctuations can be understood from the initial spatial anisotropy. The transverse momentum dependence of  $v_2$  and  $v_3$  compared to model calculations favors a small value of the shear viscosity to entropy ratio  $\eta/s$ . For the 5% most central collisions we have shown that  $v_2$  rises strongly with centrality in 1% centrality percentiles. The strong change in  $v_2$  and the small change in  $v_3$  as a function of centrality in these 1% centrality percentile classes follow the centrality dependence of the corresponding spatial anisotropies. The two-particle azimuthal correlation for the 0%–1% centrality class exhibits a double peak structure around  $\Delta\phi \sim \pi$  (the “away side”) without the subtraction of elliptic flow. We have shown that the measured anisotropic flow Fourier coefficients give a natural description of this structure.

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