

# Impact of silver ions on dielectric properties and conductivity of lithium silicate glass system mixed with red lead

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## Abstract

Glasses of the composition  $(40-x)\text{Li}_2\text{O}-10\text{Pb}_3\text{O}_4-50\text{SiO}_2$  doped with different contents of  $\text{Ag}_2\text{O}$  were fabricated and were subsequently subjected to heat-treatment. The quantitative analysis of the results of different characterization techniques indicated increasing magnitude of internal connectivity of the glass network with increasing  $\text{Ag}_2\text{O}$  content. Several dielectric parameters viz.,  $\epsilon'$ ,  $M$ ,  $Z$  and  $\sigma_{ac}$  in the frequency ( $\omega$ ) region 0.01 Hz – 1 MHz and in the region of temperature ( $T$ ) 30 °C–250 °C of these glasses were measured as functions of  $\text{Ag}_2\text{O}$  content.  $\epsilon'$  showed a decreasing tendency with  $\text{Ag}_2\text{O}$  content which may be attributed to the decrease in the space charge polarization (scp) due to the increasing order of augmentation of glass network. Such a decrease of scp was attributed to the growing presence of  $\text{PbO}_4$  units and also  $\text{Ag}^0$  particles in these samples. The plots of electric modulus vs  $\omega$  and  $T$  indicated relaxation character. Such relaxation phenomena were discussed in detail with the help of Cole-Cole diagrams. Probable dipoles responsible for such effects were recognized and analyzed.  $\sigma_{ac}$  (conductivity) exhibited a decreasing trend with  $\text{Ag}_2\text{O}$  content. The mechanism of conduction was analyzed using polaronic tunneling effect in the intermediate  $\omega$  and high  $T$  ranges, whereas in the low  $T$  range it is predicted to follow QMT model. These conclusions were observed to follow the results of various spectroscopic studies. In conclusion, presence of red lead seemed to be an obstruction for the conductivity in the fast ion conducting glasses.

**Keywords:** Red lead; Ag ions; Lithium silicate glasses; Dielectric studies; A.C. conductivity

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## 1. Introduction

Studies regarding the dielectric features of amorphous materials not only give the information about their insulating character but also give the information about their structural characteristics. In fact, one can monitor the conducting or non-conducting behaviour of glasses or glass ceramics by evaluating different dielectric parameters viz.,  $\epsilon^*$  (real and imaginary dielectric constant),  $M^*$  (the electric moduli),  $Z^*$  (impedance) and  $\sigma_{ac}$  (conductivity) over broad regions of  $\omega$  (frequency) and  $T$  (temperature) as a function of their composition. A number of such studies were reported on a variety of glass materials during the last several years [1-6].

Among various alkali silicate glasses, lithium oxide forms stable glasses with silicate at a high melting temperature. However, lithium mixed silicate glasses were reported to devitrify with greater ease than other alkali silicate glasses [7] and in fact fairly stable lithium silicate glasses can be obtained only with the mixtures containing of more than 70 mol% of  $\text{SiO}_2$ . In order to suppress devitrification, in this study lead oxide was added to the lithium silicate glasses. In the glass systems, lead oxide participates both as a network-forming oxide and as a glass network modifier. Among various forms of lead oxide, red lead ( $\text{Pb}_3\text{O}_4$ ) is an interesting heavy metal oxide. Its presence in the glass composition can cause an increase in the density and therefore an increase in the refractive index to a greater extent. As a result, the scope of applications of the lithium silicate glasses widens. In fact, such materials find their application in non-linear optical devices in addition to their well-known application as electrolytes in solid state batteries. Strictly speaking,  $\text{Pb}_3\text{O}_4$  decomposes at about  $600^\circ\text{C}$  to  $\text{PbO}$  according to the following chemical equation:  $2\text{Pb}_3\text{O}_4 \rightarrow 6\text{PbO} + \text{O}_2 \uparrow$ . During this reaction,  $\text{PbO}$  evolves from  $\text{Pb}_3\text{O}_4$  through two intermediate polymorphs,  $\alpha\text{-PbO}_2$  and  $\beta\text{-PbO}_2$ . Among these two, the  $\alpha\text{-PbO}_2$  is a relatively less stable form, whereas  $\beta\text{-PbO}_2$  is relatively

stable and exhibits good corrosion resistance even in acidic media [8,9]. The most important application of this oxide is as a cathode in the lead acid batteries [10]. However, by capturing  $2e^-$ , the  $PbO_2$  converts in to  $PbO$  as per the reaction shown above. The lead ions in  $PbO_2$  occupy mainly network forming positions with  $PbO_4$  structural units which alternate  $SiO_4$  structural units, whereas  $Pb^{2+}$  ions mainly act as glassmodifiers and also occupy network forming positions in  $PbO_4$  units [11,12].

As per the crystallographic studies, the coordination geometries of lead complexes are of two categories. One of them is holodirected. In this, the lead ligand bonds are distributed throughout the surface of an encompassing globe and are nearly tetrahedral structures. According to Cambridge Structural Database all lead (IV) structures have a holodirected coordination geometry [13]. On the other hand, lead (II) structures are hemidirected for low coordination number (2-5), in which the lead–ligand bonds are directed throughout only part of an encompassing globe, viz., there is an identifiable void in the distribution of bonds to the ligands such as pyramidal structures.

Lead(II) with electron configuration  $[Xe] 4f^{14}5d^{10}6s^2$  is one of the post-transition metal elements that exhibits the so-called “inert-pair effect” [13, 14] which means the resistance of the pair of outer electrons on Pb(II) to removal or to participation in covalent bond formation. It has been explained as a relativistic effect causing the 6s orbital to contract, thereby increasing the energy required to remove or interact with the 6s lone pair of electrons. On the other hand, the *d* and *f* orbitals are destabilized because they expand radially as a result of screening from nuclear attraction by the s and p electrons. The result is a stable, relatively inert outer lone pair of electrons in case of  $Pb^{2+}$  ions [13, 14]

Altogether, by virtue of different valence states of lead ions in  $Pb_3O_4$ , it has a strong bearing on the diffusion of  $Ag^+$  ions and  $Li^+$  ions that are responsible for conductivity in the studied glass matrix.

Among various noble metal oxides,  $\text{Ag}_2\text{O}$  dissolves easily in the amorphous matrices. During the melting of the glasses with small concentration of  $\text{Ag}_2\text{O}$ , this oxide decomposes into  $\text{Ag}^0$  (silver metallic particles) and  $\text{Ag}^+$  ions at about 420 K [15]. In general, the glasses consisting of noble metallic particles like silver particles, exhibit surface plasmon resonance (SPR) with the interaction of  $E$ -vector of the incident electromagnetic (e.m.) waves. The SPR induces high local  $E$ -field in the neighbourhood of Ag ions [16] and has strong bearing on the electrical properties of the glasses to a great degree.

Further, the addition of small concentration of  $\text{Ag}_2\text{O}$  to lithium silicate glasses makes them a super ionic conducting glass. In fact, several recent studies on electrical properties of glasses containing noble metallic particles are available in the literature [17-19]. Moreover, such materials exhibit photo-catalytic activity with oxidation stability [20]. In general, in such materials, Ag ions, and electrons ( $e^-$ s) combinedly contribute to the conductivity. For this reason, the glasses containing  $\text{Ag}_2\text{O}$  are being used as important materials for electrodes and as electrolytes in chemical cells [21]. Overall, the conductivity of these glasses is projected (i) due to polaronic hopping between Ag ions of multiple valence states and in between  $\text{Pb}^{2+}$  and  $\text{Pb}^{4+}$  ions and (ii) due to the transport of silver and lithium ions.

The content of this study includes: (i) the methods of synthesis of  $\text{Li}_2\text{O-Pb}_3\text{O}_4\text{-SiO}_2$  glasses doped with minor quantities (0 to 0.4 mol%) of  $\text{Ag}_2\text{O}$ , (ii) characterization of the glasses by means of X-ray diffraction (XRD), scanning electron microscopy (SEM), X-ray photoelectron spectroscopy (XPS) and differential thermal analysis (DTA) techniques, (iii) structural analysis of the prepared glasses as a function of  $\text{Ag}_2\text{O}$  content using IR and optical absorption spectra, (iv) detailed investigations on dielectric parameters viz., permittivity, loss tangent, electric moduli, dipolar relaxation effects, impedance diagrams and  $\sigma_a$  over a range of  $\omega$  viz.,  $10^{-2}$  to  $10^6$  Hz and in the region of temperature 30 to 250°C and (v)

analysis of these results to have an understanding over the electrical conduction process as a function of structural changes occurring in the glass matrix due to variations in Ag<sub>2</sub>O content.

## 2. Experimental

Different glasses (in mol %) in (40- $x$ )Li<sub>2</sub>O-10Pb<sub>3</sub>O<sub>4</sub>-50SiO<sub>2</sub>: $x$ Ag<sub>2</sub>O series (with  $x = 0, 0.1, 0.2, 0.3$  and  $0.4$ ) were prepared by general melt quenching technique. Based on the quantity of Ag<sub>2</sub>O the glasses were designated as A<sub>0</sub>, A<sub>1</sub>, A<sub>2</sub>, A<sub>3</sub> and A<sub>4</sub>, respectively. The appropriate amounts of reagent grade chemicals were taken in platinum crucibles and were melted at 1400 °C for 30 min. The resultant melt was then poured into brass moulds and subsequently annealed at 350 °C so as to avoid internal bubbles, cracks, etc. To ensure the entrenchment of silver particles in the samples, they were heat treated at 250 °C for about 4 h and were chilled to room temperature. Later, the samples were finely ground and optically polished. The final dimensions of the samples used for the dielectric measurements were approximately 1.0 cm x 1.0 cm x 0.1 cm. A Carl Zeiss – EvoMA15A scanning electron microscope was used to record the SEM images of the samples. Chemical make-up of the glasses was evaluated by energy dispersive spectra (EDS) recorded with an Oxford-INCA Penta FETX3. XRD patterns of these samples were recorded on A RIGAKU-D/Maxb ULTIMA III X-ray diffractometer in the range of  $2\theta$  from 15° to 65° in the steps of 0.02 °/s with a scan speed of 2 deg/min.

The thermal behaviour of glasses was investigated by differential thermal analysis (DTA) using Mettler TGA/DSC 3+ thermo balance. Measurements were performed on powder samples (cca 40 mg) placed in an open Pt crucible and heated in an oxygen atmosphere from 25 up to 1000 °C at a rate of 20 °C min<sup>-1</sup>. The results were analysed by applying the Mettler STARe 9.01 software and the glass transition temperature,  $T_g$ , was determined for each analysed sample. XPS studies were performed on PHI 5000 Versa Probe ULVAC instrument with Al K $\alpha$  (1486.6 eV) as the source of X-rays. Electron emission angle

was 45°. The stage the sample was kept on double sided carbon tape (in gifu) and the entire stage was placed into main entry chamber and kept for 12h. Etching of the samples was done with Ar<sup>+</sup> energy of 1kV for about 300 sec. 10<sup>-7</sup> Pascal was the base pressure maintained. XPS were registered with C 1s peak at 284.6 eV as the reference.

The Fourier transform-infrared (FT-IR) spectra of the samples were recorded to a resolution of 0.1 cm<sup>-1</sup> on a pre-calibrated (with a diamond crystal) Shimadzu IR TRACER 100 spectrophotometer in attenuated total reflectance (ATR) mode. A JASCO UV-VIS-NIR spectrophotometer was used to record optical absorption (OA) spectra of the samples to a precision of 0.1 nm in the wavelength region 200–2000 nm.

For measuring dielectric properties, thin gold electrodes of nearly 0.6 cm in diameter, were sputtered on both sides of the glasses with SC7620 sputter coater. The dielectric and a.c. conductivity measurements were carried out on Novocontrol Alpha-AN Dielectric Spectrometer in the range of frequency 0.01 Hz- 1 MHz and in the range of temperature 30 - 250 °C (up to a precession of ±0.2°C). Equivalent circuits modelling and CNLLSQ (Complex Nonlinear Least Squares fitting) process was adopted for analysing the complex impedance diagrams.

### 3. Results and Discussion

Table 1 summarizes the concentration of Ag<sup>+</sup> ions and their average distance in the glass samples evaluated using the measured values of density and the average molecular weight.

Images of Li<sub>2</sub>O-Pb<sub>3</sub>O<sub>4</sub>-SiO<sub>2</sub> glasses containing varied amounts of Ag<sub>2</sub>O were presented in Fig. 1. The samples with low concentrations of Ag<sub>2</sub>O were found to be colourless and appeared to be light-yellow to brownish-yellow in colour as the Ag<sub>2</sub>O content was increased. This demonstrated formation of forming colloidal Ag<sup>0</sup> particles (formed by

reducing  $\text{Ag}^+$  ions) and suggested larger concentration of Ag nanoparticles (Ag NPs) in the samples containing relatively larger concentration of silver oxide [22].

SEM images of post-heated glasses with varied contents of  $\text{Ag}_2\text{O}$  were shown in Fig. 2 which indicated the presence of well-defined unevenly dispersed crystallites (with a few  $\mu\text{m}$  size) embedded in the remnant glass phase. As the content of  $\text{Ag}_2\text{O}$  was increased, an increasing fraction of crystal grains with growing coalescence was observed. The chemical composition of the glasses was verified by EDS; the spectra suggested intact of all the expected elements in these glasses. Such a spectrum for the A<sub>4</sub> glass was shown in Fig. 2,

X-ray diffractograms (Fig.3) of the heat-treated glasses showed peaks corresponding to  $\text{Ag}_2\text{O}$  (JPCDS 19-1155),  $\text{AgO}$  (JPCDS 43-1038) crystalline phases and  $\text{Ag}^0$  (JPCDS 04-0783) particles [23-25]. Among these,  $\text{Ag}_2\text{O}$  is a cubic phase, whereas  $\text{AgO}$  is a monoclinic ( $a = 5.85$ ,  $b = 3.47$ ,  $c = 5.49$ ,  $\beta = 107.5^\circ$  with space group  $\text{P2}_1/\text{c}$ ) as shown in this figure [26]. With an increase in the  $\text{Ag}_2\text{O}$  content, a gradual growth of the diffraction peaks related to  $\text{Ag}^0$  particles was observed. These observations suggested a growing concentration of  $\text{Ag}^0$  particles in the heat-treated samples.

DTA traces of the investigated glasses were shown in Fig. 4. The traces exhibited a comparatively weak endothermic effect due to glass transition in the temperature range 400–415°C. The thermograms exhibited multiple-exothermic peaks due to crystallization in the temperature range 500-700 °C with a prominent peak in the temperature region 500-600 °C. Another significant endothermic change due to re-melting of the samples was observed at about 900 °C. A brief summary of the results thermal analysis of these samples was given in Table 2. The results of differential thermal analysis suggested that the heat-treated samples comprised of multiple crystal phases. Further, a decrease in  $T_{\text{cl}} - T_{\text{g}}$  was observed with increase in  $\text{Ag}_2\text{O}$  content (Table 2). The observed decrease suggested an increased degree of devitrification of the glass samples.

Fig. 5, shows XPS spectrum of the glass A<sub>2</sub>. The deconvoluted and fitted spectrum exhibited three peaks at 137.5 eV, 137.9 eV and 138.4 eV corresponding to PbO<sub>2</sub>, PbO and Pb<sub>3</sub>O<sub>4</sub> region of Pb 4f<sub>7/2</sub>. Three other peaks at 142.36 eV, 142.76 eV and 143.26 eV corresponding to PbO<sub>2</sub>, PbO and Pb<sub>3</sub>O<sub>4</sub> region of Pb 4f<sub>5/2</sub> were also observed in this spectrum [27]. With an increase in Ag<sub>2</sub>O concentration, the peak corresponding to PbO<sub>2</sub> exhibited a gradual growth at the expense of Pb<sub>3</sub>O<sub>4</sub> peak (inset of Fig. 5). This observation suggested that there was a gradual growth of Pb<sup>(IV)</sup>-O bondings in the glass network that facilitated the increased degree of structural augmentation. Such augmentation was predicted to be a hindrance for the easy transport of charge carriers that contribute to the conductivity

XPS spectrum of the glass A<sub>2</sub> recorded in the energy region of Ag 3d was presented in Fig. 6. The spectrum showed the peaks at 367.2 and 373.1 eV, respectively, associated with Ag<sup>+</sup> (Ag 3d<sub>3/2</sub>) and Ag<sup>0</sup> (Ag 3d<sub>5/2</sub>) NPS [28]. A feeble kink corresponding to Ag<sup>2+</sup> ions was also observed in this spectrum at 367.5 eV [28, 29]. In the inset of this diagram variations of the intensities of the peaks corresponding to Ag<sup>+</sup> ions and Ag<sup>0</sup> metallic particles were plotted against concentration of Ag<sub>2</sub>O. The inset showed the gradual growth of the peak corresponding to Ag<sup>0</sup> metallic particles, whereas the peak related to Ag<sup>+</sup> ions that contribute to the ionic conductivity gradually decreased.

The optical absorption spectra of Li<sub>2</sub>O–Pb<sub>2</sub>O<sub>3</sub>–SiO<sub>2</sub>: Ag<sub>2</sub>O glass recorded in the wavelength region 250–800 nm were presented in Fig. 7. Pure glass (A<sub>0</sub>) exhibited cut-off wavelength at 330.5 nm. With the gradual increase in Ag<sub>2</sub>O from 0.1 to 0.4 mol%, the edge exhibited spectrally blue shift. Moreover, these spectra exhibited an absorption band in the spectral region 340–460 nm. In general Ag<sup>+</sup> ions get transformed to metastable Ag<sup>0</sup> (4d<sup>10</sup>5s<sup>1</sup>) metallic particles [30]. These unstable Ag<sup>0</sup> particles interact with Ag<sup>+</sup> ions and Ag<sup>+</sup>-Ag<sup>0</sup> clusters form in the material. There occurs surface plasmon resonance (SPR) between the coherent oscillations of e<sup>-</sup>s on the top of Ag<sup>+</sup>-Ag<sup>0</sup> groups and the e.m. waves incident upon



the material. The visualised peak in the spectra represented such resonance [31]. Further, this band exhibited a growth with  $\text{Ag}_2\text{O}$  content. It indicated increasing presence of  $\text{Ag}^0$  particles with the rise of  $\text{Ag}_2\text{O}$  content in the glasses. Additionally, the SPR band exhibited red shift with the concentration of  $\text{Ag}_2\text{O}$ . Some of the reasons for that red shift include the increasing dimensions of the clusters of silver particles, scattering contribution to SPR band and there may be change in the absorption from dipoles to multipoles etc. [32-34].

To estimate the optical band gap ( $E_o$ ) and to study its dependence on  $\text{Ag}_2\text{O}$  concentration, Tauc plots ( $(\alpha\hbar\omega)^{1/2}$  vs  $\hbar\omega$ ) were drawn in Fig. 8. The value of  $E_o$  (evaluated from the extrapolation of the curves to the X-axis) exhibited an increasing trend with  $\text{Ag}_2\text{O}$  concentration (Table 3). XPS results indicated that there was an increase in  $\text{Pb}^{4+}$  ion concentration with an increase in  $\text{Ag}_2\text{O}$  content. Tetravalent lead ions take part in glass network forming positions through  $\text{Pb}^{\text{IV}}\text{-O-Si}^{\text{IV}}$  linkages. Such interconnectivity among various structural groups caused a decrease in the degree of localization of free  $e^-$ s and decreased the concentration of donor centres and lead to enlarged  $E_o$ .

Yet, a variation in  $E_o$  with  $\text{Ag}_2\text{O}$  content is connected with the interactions of  $p$ - $e^-$ s of Pb and  $d$ - $e^-$ s of Ag ions. XPS and even XRD results revealed an increase in  $\text{Ag}^+$  ions' concentration which were reduced to  $\text{Ag}^0$  particles. This caused a decrease in the donor centres' concentration. For this reason, the electrons trapped at energy states of  $\text{Ag}^+$  ions get distanced to a larger degree with the empty  $4d$  energy levels of  $\text{Ag}^0$  particles. As a result, width of the band constitutes energy levels of  $\text{Ag}^0$  impurity contracted in the band gap of the host glass and caused to enlarge  $E_o$  with the increase of  $\text{Ag}_2\text{O}$  content.

Shifting of SPR band towards longer wavelength with increased  $\text{Ag}_2\text{O}$  content (Table 3) indicated mutual exchange of polarons among  $\text{Ag}^0$ - $\text{Ag}^+$  pairs [35]. Effective radial wave-function of  $e^-$ s trapped at less deeper traps in the band gap was shorter because of insignificant local oscillations of potentials compared to transfer integral  $J$ . As a result, the

magnitude of superposition of wavefunctions of the electrons in the closely spaced sites reduced and facilitated the creation of polaron. In summary, the results of optical absorption of Ag<sub>2</sub>O doped Li<sub>2</sub>O-Pb<sub>3</sub>O<sub>4</sub>-SiO<sub>2</sub> glasses suggested the possible polaronic transfer between Ag<sup>0</sup>-Ag<sup>+</sup> pairs. Such type of transfer predicted to be vital for analysing the dielectric results [36].

The IR spectra of the studied glasses (Fig. 9) exhibited two prominent bands ascribed to asymmetric and symmetric stretching vibrations O-Si-O of SiO<sub>4</sub> units at about 1012 cm<sup>-1</sup> and at 790 cm<sup>-1</sup>, respectively [37]. Additionally, the spectra also exhibited two vibrational bands at 578 cm<sup>-1</sup> (identified as the characteristic vibrations of Pb<sup>(IV)</sup>-O bond of PbO<sub>4</sub> units [8]) and at 460 cm<sup>-1</sup> (conventional band due to Pb<sup>(II)</sup>-O vibrations [37]). With the gradual increase of Ag<sub>2</sub>O concentration, the bands due to symmetrical vibrations of O-Si-O groups and also those of PbO<sub>4</sub> structural groups were observed to grow. Such growth indicated a growing degree of structural augmentation. Positions of different vibrational bands visualized in these spectra were summarized in Table 4.

Variation of dielectric constant  $\epsilon'(\omega)$  with  $\omega$  at different temperatures (for two of the glasses viz., A<sub>1</sub> and A<sub>2</sub> were shown in Fig. 10 (a) and with  $T$  at various  $\omega$  (for A<sub>3</sub> and A<sub>4</sub> glasses) were shown in Fig. 10 (b). With the increase in  $\omega$ ,  $\epsilon'(\omega)$  exhibited a decreasing trend and a high frequency plateau value  $\epsilon'_{\infty}(\omega)$  was reached. With temperature,  $\epsilon'(\omega)$  exhibited a larger dispersion at high  $T$  and larger values at low frequencies (increased from 10<sup>2</sup> to 10<sup>5</sup> with increase of temperature from 30 to 250 °C). Among different electrical polarizations, the space charge polarization (piling up of charge carriers at the electrodes) caused to increase the value of  $\epsilon'(\omega)$  at lower frequencies. In general, the quantum of this polarization is dependent on long range diffusion (which depends on the degree of structural defragmentation) of free charge carriers). To put it differently, the larger increase of  $\epsilon'(\omega)$  at low  $\omega$  region was because of easy diffusion of conducting species in the material. In high  $\omega$

region, the oscillation frequency of the conducting species was relatively lower. For this reason,  $\varepsilon'(\omega)$  reached high frequency plateau  $\varepsilon'_\infty(\omega)$  as shown in Fig.10(a).

In the inset of Fig.10(a), the dependence of  $\varepsilon'(\omega)$  on the concentration of  $\text{Ag}_2\text{O}$  (at a fixed  $\omega$  and  $T$ ) was plotted.  $\varepsilon'(\omega)$  showed a descending tendency with the increase of  $\text{Ag}_2\text{O}$  concentration. XPS and IR spectral studies suggested increased proportion of  $\text{Pb}^{4+}$  ions which occupy tetrahedral positions in this glass network and cross-link with  $\text{SiO}_4$  units. Such cross-linkages obstructed the passage of charge carriers towards electrodes [38, 39] and decreased the value of dielectric constant.

Variations of loss tangent ( $\tan \delta$ ) with frequency and temperature for the sample A4 were presented in Fig. 11. The variations indicated dipolar relaxation behaviour more clearly in the low  $\omega$  range. Nevertheless, in the high  $\omega$  region the dipolar effects were not distinct; electrode phenomenon appeared to be dominating over the inherent dipolar effects in this frequency range. Hence, to throw more light on the inherent dipolar effects, electric moduli formulation was adopted. According to this formalism, electric modulus  $M^*$ , is given by

$$M^* = M' + iM'', \quad (1)$$

$$\text{Where, } M' = \frac{\varepsilon'(\omega)}{(\varepsilon'(\omega))^2 + (\varepsilon''(\omega))^2} \quad (2)$$

$$\text{And } M'' = \frac{\varepsilon''(\omega)}{(\varepsilon'(\omega))^2 + (\varepsilon''(\omega))^2}; \quad (3)$$

Using the values of  $\varepsilon'(\omega)$  and  $\varepsilon''(\omega)$ ,  $M'$  and  $M''$  were evaluated and their variations with  $\omega$  and  $T$  were presented Fig. 12. A clear dipolar relaxation character of the inherent dipoles could be visualized from these plots.

Using the resonant frequency  $\omega_r$  (points of intersection of  $M'$  and  $M''$ ), the relaxation time  $\tau$  was determined and its variation with  $T$  was plotted in Fig. 13 for all the samples. The variation showed a decreasing tendency with  $T$  and found to be independent of temperature

at highertemperatures. Variation of  $\tau$  with  $\text{Ag}_2\text{O}$  content at a fixed  $T$  was observed to increase (Table 5). From the plots of  $\ln \tau$  vs  $1/T$ , the activation energy  $W_d$  for dipoles was estimated using the equation  $\tau = \tau_0 \exp (W_d/k_B T)$  and presented in Table5.  $W_d$  exhibited an increase with the increase in the concentration of  $\text{Ag}_2\text{O}$ . Such increase of  $W_d$  and also  $\tau$  with the content of  $\text{Ag}_2\text{O}$  suggested decreased magnitude of freedom for dipoles to orient in the field direction.

To shed additional light on dipolar effects of the samples, Cole-Cole diagrams between  $M''$  vs  $M'$  at higher temperatures were drawn and presented in Fig. 14 for the glass samples A<sub>3</sub> and A<sub>4</sub>. The plots were observed to be semi-circular arcs with centers lying below the X-axis. The obtained values of angle ( $\alpha'$ ) between line connecting the center of arc and the X-axis at the origin are presented in Table5.  $\alpha'$  was found to decrease with increase of  $\text{Ag}_2\text{O}$  content and suggested decreasing degree of distribution of  $\tau$ . In general, spreading of relaxation times is due to multiple types of dipoles with various dipole moments [40]. Additionally, entrenchment of dipoles in asymmetrical potential regions also accounts for such spreading[7]. The decreased value of  $\alpha'$  suggested the decreasing asymmetric environment for the dipoles as was also suggested by IR spectral studies. The complexes of divalent silver ions and also divalent lead ions did contribute for dipolar relaxation effects. The growing concentration of tetravalent lead ions that occupy the glass network forming positions resist the freedom for dipoles to oscillate and increased activation energy for dipoles.

Figs. 15(a) and (b), respectively, represent the variations of real ( $Z'$ ) and imaginary ( $Z''$ ) parts of impedance with frequency at different temperatures for the glass A<sub>2</sub>. Out of the two components of impedance, the real component  $Z'$  (a pure resistance component) showed a decrease with  $\omega$  (even though it exhibited a minor non-linearity in the low  $\omega$  range, (likely due to electrode effect)). The imaginary part  $Z''$ , however, showed large spreading with  $T$  in

the low frequency region ( $< 1$  kHz). Moreover, nearly up to 1 kHz, this component exhibited a linear increase with  $\omega$  especially at higher temperatures. This type of behaviour indicated inductive reactance component ( $L\omega$ ). However, at higher frequencies (beyond 1kHz) the  $Z''$  component with frequency exhibited an inverse proportionality. Such behaviour is a characteristic of capacitive reactance (*viz.*,  $1/C\omega$ ). To throw more light on the physical significance of these diagrams, in Fig. 15(c),  $Z''$  and  $M''$  vs frequency at 200°C were plotted for A<sub>2</sub> sample. The peaks of these two quantities were found to be slightly deviated and suggested that long-range interaction was the main reason for such a confined relaxation character of these glasses [41].

To know the dependence of magnitude of the impedance of Li<sub>2</sub>O-Pb<sub>3</sub>O<sub>4</sub>-SiO<sub>2</sub>:Ag<sub>2</sub>O glass samples on the content of Ag<sub>2</sub>O, impedance plots (or Nyquist plots) between  $Z''$  vs  $Z'$  were drawn. Figs. 16 (a), (b) and (c) represented such plots for the sample A<sub>2</sub> drawn in different temperature regions. Area enclosed by the curve exhibited a decrease with temperature and indicated a decrease in the impedance of the samples with temperature. The comparison of impedance plots for all the samples drawn at 250 °C was shown in Fig. 16(d) along with the equivalent circuit diagram. The area under the curves was observed to increase with increase of Ag<sub>2</sub>O content and suggested an increase of  $Z$ . Moreover, the diagrams exhibited semi-circular shape with inclined spur (at low  $\omega$  and high  $T$  regions) due to thermally stimulated mobility of the conducting species.

These diagrams further indicated two kinds of conducting species in the glass samples. Semi-circular arc associated with bulk conduction and the inclined spur represented the piling of charge carriers on surface electrodes of the glasses. Larger magnitude of such spur, suggested larger concentration of conducting species that accumulated at the electrodes. In this investigation, a descending size of magnitude of this spur along with growth of the area enclosed by the semicircle was observed with the increase of Ag<sub>2</sub>O concentration. This

behaviour indicated increasing magnitude of the resistance for the migration of  $\text{Li}^+$  and  $\text{Ag}^+$  ions (conducting ions) towards the electrodes. The impedance behaviour was analysed further, by Bode plots (impedance  $|Z|$  and  $\tan\phi = Z''/Z'$  vs  $\omega$ ). Such diagram for a particular glass A<sub>4</sub> drawn at 200 °C was presented in Fig. 16(e). The graph suggested that  $|Z|$  was higher at low  $\omega$  region and lower in the high  $\omega$  region and indicated that electronic/polaronic (or Faradic) current was responsible for the impedance [42, 43].

In Figs. 17 (a) and 17 (b),  $\ln\sigma_{ac}$  (ac conductivity) was plotted against  $\omega$  and against  $1/T$  for the glass A<sub>4</sub>, respectively. In the high  $T$  region and in the low  $\omega$  range ( $\sigma_{ac}(\omega) \rightarrow \sigma_{dc}$ ) the plots were found to be near frequency independent; it was observed to be near temperature invariant on  $T$  in the high  $\omega$  and low temperature regions and it varied linearly with  $\omega$  viz.,  $\sigma_{ac}(\omega) \propto \omega$ . The general relationship between  $\sigma_{ac}(\omega)$  and  $\omega$  however, seemed to be  $\sigma_{ac}(\omega) \propto \omega^s$ , where  $s=1$  in the high  $\omega$  and low  $T$  regions; the  $\omega$  exponent  $s$  is the universal dielectric response (UDR) exponent; normally,  $s < 1$ .

In the intermediate  $\omega$  region, a ‘*to* and *fro*’ hopping of conducting species (polarons in this case) occurred due to coulombic repulsion. Structural investigations suggested that Pb and Ag ions existed in multiple valence states. Hence, the polaronic exchange between  $\text{Pb}^{2+}$  and  $\text{Pb}^{4+}$  and  $\text{Ag}^0$  to  $\text{Ag}^+$  ions took place and contributed to the  $\sigma_{ac}$  in addition to the ionic contribution from  $\text{Li}^+$  and  $\text{Ag}^+$  ions. In the inset of Fig. 17 (a) the dependence of  $\sigma_{ac}$  (measured at 1 kHz and 250 °C) on the concentration of  $\text{Ag}_2\text{O}$  was plotted. The plot suggested a gradual decrease of  $\sigma_{ac}$  with  $\text{Ag}_2\text{O}$  content. The obtained values of  $\sigma_{ac}$  were found to be comparable with that of various other glass systems containing  $\text{Ag}_2\text{O}$  (Table 6) [28, 44-47]

From the plots of  $\ln\sigma_{ac}$  vs  $1/T$ , the activation energy  $W_{ac}$  was evaluated and it was found to increase with increase of  $\text{Ag}_2\text{O}$  concentration (Table 5). Such increment suggested a

decreasing magnitude of free volume space available for easy movement of conducting species in the glasses with increase of Ag<sub>2</sub>O content.

UDR exponent  $s$  was estimated from the plots of  $\log \sigma_{ac}(\omega)$  vs  $\log \omega$  (drawn in the intermediate  $\omega$  region). Such plot was presented in Fig. 18 for the glass A<sub>4</sub>. The value of  $s$  showed a decreasing trend with  $T$  (inset of Fig. 18). The decrease of the exponent suggested that conduction consists of polaronic tunneling[44]. Variation of  $s$  with Ag<sub>2</sub>O content indicated an increasing trend (Table 5). Such increase indicated a decrease in the magnitude of free space available for the diffusion of Li<sup>+</sup> and Ag<sup>+</sup> ions with increase of Ag<sub>2</sub>O content [48]. From the analysis of the results of structural studies, we found that in the samples containing larger concentration Ag<sub>2</sub>O, there was a larger concentration of Pb<sup>IV</sup>-O units that cross link with SiO<sub>4</sub> units. Such cross linkages obstruct the easy transport of conducting species like Li<sup>+</sup> and Ag<sup>+</sup> ions. As a result, a decrement in the conductivity was observed.

Further, a near invariance of  $\sigma_{ac}$  with  $T$  in the low  $T$  region especially in the high frequency range (Fig. 17(b)) was observed. Such behavior of conductivity was analyzed using quantum mechanical tunneling (QMT) model [49]. As per this model, the density of defect energy states  $N(E_F)$  was evaluated using the expression

$$N(E_F) = \left[ \frac{3\sigma(\omega)}{\pi e^2 k_B T (\alpha'')^{-5} \omega \left[ \ln \frac{\nu_{ph}}{\omega} \right]^4} \right]^{1/2} \quad (4)$$

In Eq. (4),  $\alpha''$  is the electronic wave function decay constant (evaluated to be  $\sim 0.45(\text{\AA})^{-1}$  by plotting  $\log \sigma_{ac}$  vs  $R_i$ , inter-ionic separation of Ag ions),  $\nu_{ph}$  is phonon frequency,  $e$  is the electron charge. Value of  $N(E_F)$  evaluated (at 1 kHz and at 473 K) showed a decreasing trend with Ag<sub>2</sub>O concentration (Table 5). This result further confirmed that there was a decreasing concentration of defects in the glass (that caused to lower the  $\sigma_{ac}$ ) with increase of Ag<sub>2</sub>O concentration.

## 4 Conclusions

$\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2:\text{Ag}_2\text{O}$  glasses were synthesized. The results of characterization suggested that the glasses were composed of  $\text{Ag}^+$ ,  $\text{Ag}^{2+}$ ,  $\text{Pb}^{4+}$ ,  $\text{Pb}^{2+}$  ions and  $\text{Ag}^0$  metallic particles. The results further suggested an increase in the concentration of tetrahedral  $\text{Pb}^{4+}$  ions which cross-link with  $\text{SiO}_4$  units. The observed decrease in the dielectric permittivity with  $\text{Ag}_2\text{O}$  concentration was attributed to the decreased magnitude of scp because of decreased concentration of free volume imperfections. The observed dipolar relaxation effects were analysed using Cole-Cole diagrams. Ac conductivity ( $\sigma_{ac}$ ) was ascribed to the polaronic transfer between multiple valence states of Ag ions and also Pb ions in addition to the diffusion of  $\text{Li}^+$  and  $\text{Ag}^+$  ions. The observed decrease of  $\sigma_{ac}$  with  $\text{Ag}_2\text{O}$  content was attributed to the increased degree of polymerization of the glass network. Overall, it appeared the presence of red lead was an obstruction for the conductivity even in fast ion conducting glasses.

### Credit author statement

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### Caption for Tables

**Table 1** Details of composition (all in mol%) of the glass samples used.

**Table 2** Summary of the data on differential thermogravimetric studies

**Table 3** Summary of the data on optical absorption spectra of  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2:\text{Ag}_2\text{O}$  glasses. The band positions are measured to an accuracy of  $\pm 1$  nm.

**Table 4** Summary of the data of band positions (in  $\text{cm}^{-1}$ ) of various structural units in FT-IR spectra of  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2:\text{Ag}_2\text{O}$  glasses. The values of the peak positions are measured to precision of  $\pm 1$   $\text{cm}^{-1}$ .

**Table 5** Summary of the data on a.c. conductivity and other parameters of  $\text{Ag}_2\text{O}$  doped  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses.

**Table 6** A.c. conductivity of various glass systems doped with  $\text{Ag}_2\text{O}$  measured at 1 kHz and 200 °C)

### Caption for Figures

**Fig. 1** Photographs of the heat treated  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

**Fig. 2** SEM images of heat treated  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ . The figure also consists of EDS for the sample A<sub>4</sub>.

**Fig. 3** XRD profiles of heat treated  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

**Fig. 4** DTA traces of heat treated  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

**Fig. 5** Deconvoluted and fitted XPS spectra of PbO, PbO<sub>2</sub> and Pb<sub>3</sub>O<sub>4</sub> in the glass sample A<sub>2</sub> and inset represents the intensity variation of PbO<sub>2</sub> and Pb<sub>3</sub>O<sub>4</sub> (4f<sub>7/2</sub>) with the concentration of silver oxide.

**Fig. 6** Deconvoluted and fitted XPS spectra of Ag<sup>0</sup> and Ag<sup>+</sup> ions in the glass samples A<sub>2</sub> and the insets represents the intensity variation of Ag<sup>0</sup> and Ag<sup>+</sup> ions with the concentration of silver oxide.

**Fig. 7** Optical absorption spectra of  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$  recorded after heat treatment.

**Fig. 8** Tauc plots of heat treated  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

**Fig. 9** IR spectra of heat treated  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

**Fig. 10 (a)** Variations of dielectric constant with frequency measured at different temperatures for the sample A<sub>1</sub> and A<sub>2</sub>.

**Fig. 10 (b)** Variations of dielectric constant with temperature measured at different frequencies for the sample A<sub>3</sub> and A<sub>4</sub>.

**Fig. 11** Variations of dielectric loss (a) with frequency measured at different temperatures and (b) with temperatures for the sample A<sub>4</sub>

**Fig. 12** Variations of electric moduli M' and M'' (a) with frequency at different temperatures and (b) with temperature at different frequencies for the sample A<sub>3</sub>

**Fig. 13** The plots of relaxation times vs temperature for Li<sub>2</sub>O-Pb<sub>3</sub>O<sub>4</sub>-SiO<sub>2</sub> glasses doped with different concentrations of Ag<sub>2</sub>O. In this figure the variations of  $\tau$  evaluated at 470 K with the concentration of Ag<sub>2</sub>O is also presented.

**Fig. 14** Cole-Cole plots for the samples A<sub>3</sub> and A<sub>4</sub>

**Fig. 15** Variations of Z' vs Z'' with frequency at different temperatures for the samples A<sub>2</sub>

**Fig. 15(c)** Variation of M'' ( $\omega$ ) and Z''( $\omega$ ) with frequency at 200 °C for the glass A<sub>2</sub>.

**Fig. 16 (a)** The plots of Z' vs Z'' drawn in the temperature regions (a) 30-50 °C (b) 100-150 °C (c) 225-250 °C for the sample A<sub>4</sub>.

**Fig. 16(d)** The comparison plots of impedance of temperature of Li<sub>2</sub>O-Pb<sub>3</sub>O<sub>4</sub>-SiO<sub>2</sub> glasses doped with different concentrations of Ag<sub>2</sub>O drawn at 250 °C. In the equivalent circuit diagram the parallel CR circuit accounts for semi-circular path of the diagram whereas additional series capacitor represents the inclined spur part.

**Fig. 16(e)** Bode plots (plots of total impedance |Z| and the phase angle,  $\tan\phi = Z''/Z'$  vs  $\omega$ ) for a selected glass sample A<sub>4</sub> drawn at 200 °C

**Fig. 17** The plots of  $\sigma_{ac}$  (a) vs frequency drawn at different temperatures and (b) vs 1/T drawn at different frequencies for the glass samples A<sub>4</sub>. Inset represents the variation of  $\sigma_{ac}$  with concentration of Ag<sub>2</sub>O measured at 250°C and 1 kHz.

**Fig. 18** Plots of  $\log \sigma_{ac}(\omega)$  versus  $\log \omega$  for the glass sample A<sub>2</sub> drawn at different temperatures and to evaluate frequency exponents.

**Table 1**

Details of composition (all in mol%) of the glass samples used.

Glass	Conc. Ag <sub>2</sub> O (mol%)	Density <i>d</i> (g/cm <sup>3</sup> ) (±0.001)	Silver ion conc. N <sub>i</sub> (x10 <sup>19</sup> /cm <sup>3</sup> ) (±0.01)	Interionic distance of silver ions R <sub>i</sub> (Å)(±0.01)	Polaron radius R <sub>p</sub> (Å) (±0.01)
A <sub>0</sub>	0	3.998	--	--	--
A <sub>1</sub>	0.1	4.027	1.366	4.18	1.68
A <sub>2</sub>	0.2	4.031	2.731	3.32	1.33
A <sub>3</sub>	0.3	4.038	4.101	2.89	1.16
A <sub>4</sub>	0.4	4.083	5.466	2.63	1.06

**Table 2**

Summary of the data on differential thermo gravimetric studies

Glass	Glass transition temperature $T_g (\pm 1 \text{ } ^\circ\text{C})$	Crystallization temperature $T_c (\pm 1 \text{ } ^\circ\text{C})$	$T_c - T_g (\pm 1 \text{ } ^\circ\text{C})$
A <sub>0</sub>	404	591	187
A <sub>1</sub>	413	584	171
A <sub>2</sub>	409	574	165
A <sub>3</sub>	407	565	158
A <sub>4</sub>	405	549	144



**Table 3**

Summary of the data on optical absorption spectra of  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2:\text{Ag}_2\text{O}$  glasses. The band positions are measured to an accuracy of  $\pm 1$  nm.

Absorption Band details	Glass	A <sub>0</sub>	A <sub>1</sub>	A <sub>2</sub>	A <sub>3</sub>	A <sub>4</sub>
SPR band position		--	406	409	411	419
Cut-off wavelength (nm)		330	312	300	291	276
Optical bandgap $E_o$ (eV)		3.31	3.50	3.59	3.76	3.80

**Table 4**

Summary of the data of band positions (in  $\text{cm}^{-1}$ ) of various structural units in FT-IR spectra of  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2:\text{Ag}_2\text{O}$  glasses. The values of the peak positions are measured to a precision of  $\pm 1$   $\text{cm}^{-1}$ .

Glass	Asymmetric stretching vibrations of O-Si-O bonds	Symmetric stretching vibrations of $\text{SiO}_4$ units	$\text{Pb}^{(\text{IV})}\text{-O}$ bond	$\text{PbO}_4$ units
A <sub>0</sub>	992	774	578	461
A <sub>1</sub>	1016	772	574	458
A <sub>2</sub>	1026	769	568	454
A <sub>3</sub>	1032	767	563	452
A <sub>4</sub>	1036	764	556	449

**Table 5**

Summary of the data on a.c. conductivity and other parameters of Ag<sub>2</sub>O doped Li<sub>2</sub>O–Pb<sub>3</sub>O<sub>4</sub>–SiO<sub>2</sub> glasses.

Glass	Exponent $s$ ( $\pm 0.01$ )	$N(E_F)$ ( $\times 10^{21}$ , $\text{eV}^{-1}/\text{cm}^3$ ) ( $\pm 0.01$ )	A.E. for conduction $W_{ac}$ (eV) ( $\pm 0.01$ )	A.E. for dipoles (eV) ( $\pm 0.01$ )	$\tau_M$ ( $\mu\text{s}$ ) (at $T = 190^\circ\text{C}$ ) ( $\pm 0.01$ )	Spreading factor (rads) $\alpha'$ ( $\pm 0.01$ )
A <sub>0</sub>	0.49	2.15	0.69	0.32	86.4	0.59
A <sub>1</sub>	0.53	2.06	0.74	0.36	130	0.52
A <sub>2</sub>	0.58	1.92	0.80	0.41	152	0.45
A <sub>3</sub>	0.64	1.83	0.87	0.45	191	0.42
A <sub>4</sub>	0.69	1.68	0.91	0.53	196	0.36

**Table 6**

A.c. conductivity of various glass systems doped with Ag<sub>2</sub>O measured at 1 kHz and 200 °C

S.No.	Glass system	A,c, conductivity (ohm-cm) <sup>-1</sup>	
1	Ag <sub>2</sub> O doped Li <sub>2</sub> O–Nb <sub>2</sub> O <sub>5</sub> –P <sub>2</sub> O <sub>5</sub> glasses	$8.2 \times 10^{-8}$	Ref. [28]
2	Ag <sub>2</sub> O doped sodium antimonate glass ceramics	$7.5 \times 10^{-8}$	Ref. [44]
3	Ag <sub>2</sub> O doped LiF–V <sub>2</sub> O <sub>5</sub> –P <sub>2</sub> O <sub>5</sub> glass	$8.2 \times 10^{-8}$	Ref. [45]
4	AgI–Ag <sub>2</sub> O–V <sub>2</sub> O <sub>5</sub> glasses	$6.84 \times 10^{-6}$	Ref. [46]
5	AgI doped silver phosphate glass	$5.0 \times 10^{-7}$	Ref. [47]
6	Li <sub>2</sub> O–Pb <sub>3</sub> O <sub>4</sub> –SiO <sub>2</sub> : Ag <sub>2</sub> O (A <sub>4</sub> )	$8.64 \times 10^{-8}$	Present work

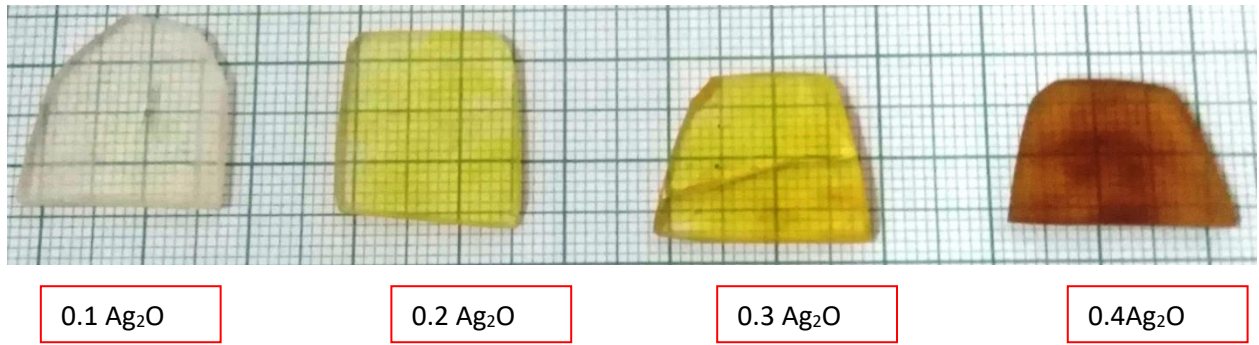


Fig. 1 Photographs of the heat treated  $\text{Li}_2\text{O-Pb}_3\text{O}_4\text{-SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

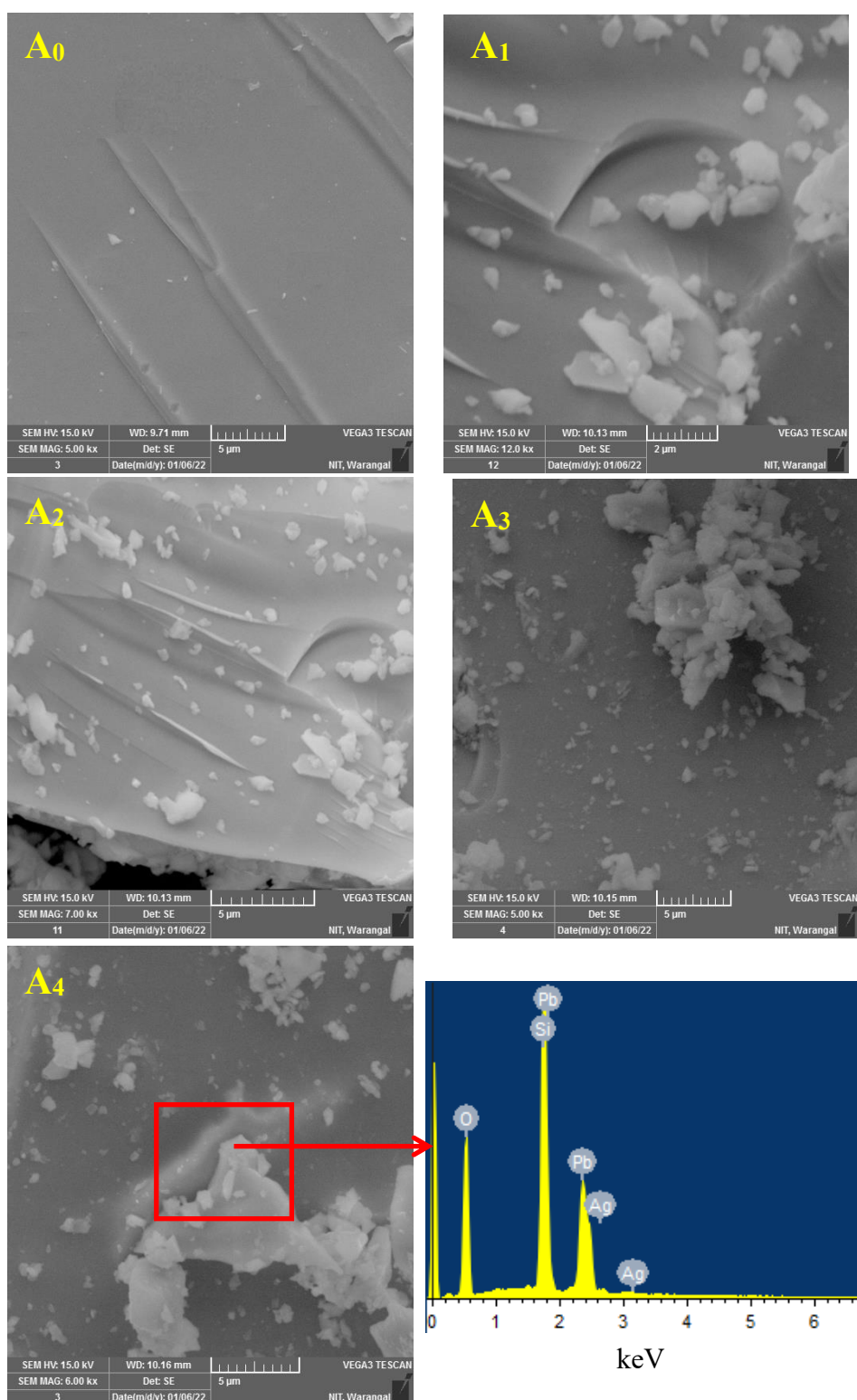


Fig. 2 SEM images of heat treated  $\text{Li}_2\text{O}-\text{Pb}_3\text{O}_4-\text{SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ . The figure also consists of EDS for the sample A4.

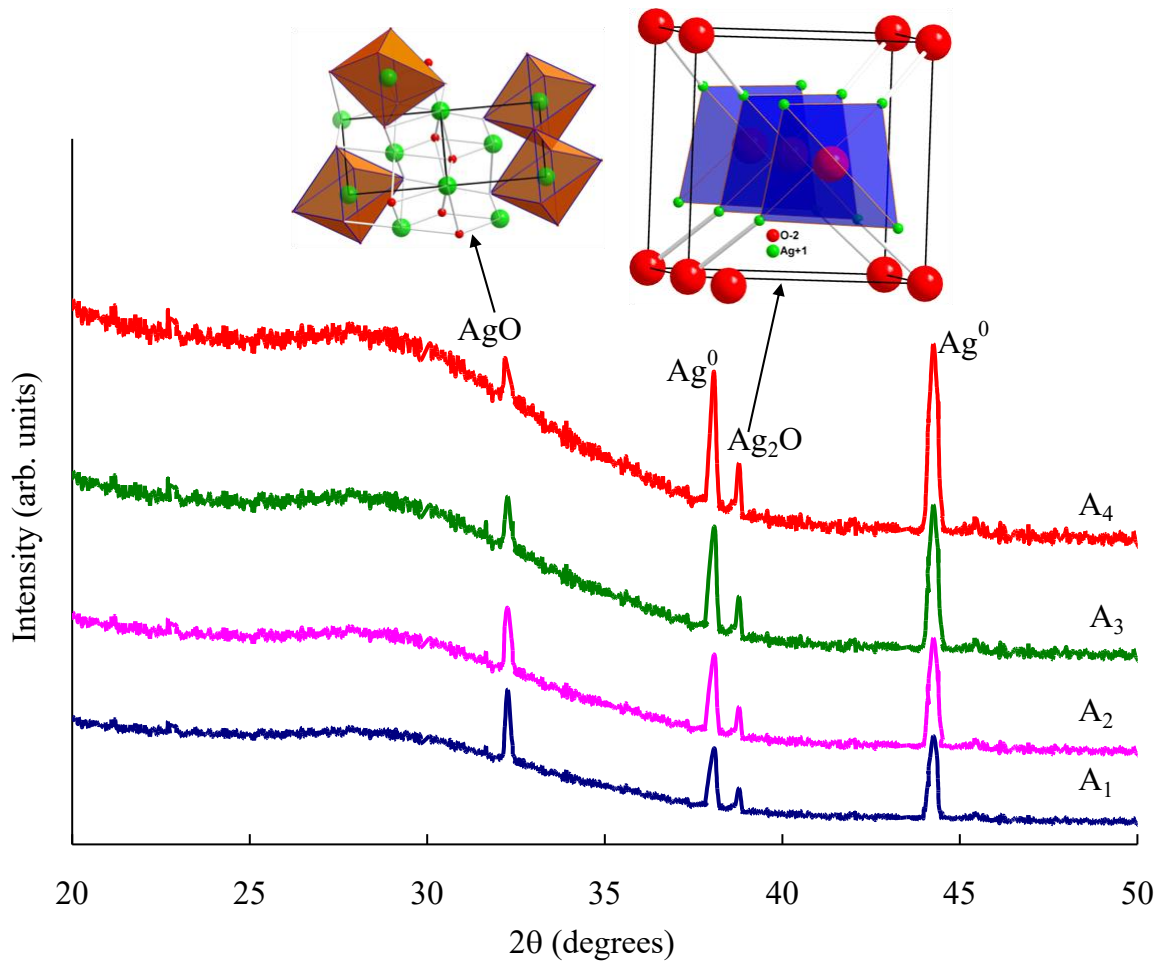


Fig. 3 XRD profiles of heat treated  $\text{Li}_2\text{O-Pb}_3\text{O}_4\text{-SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

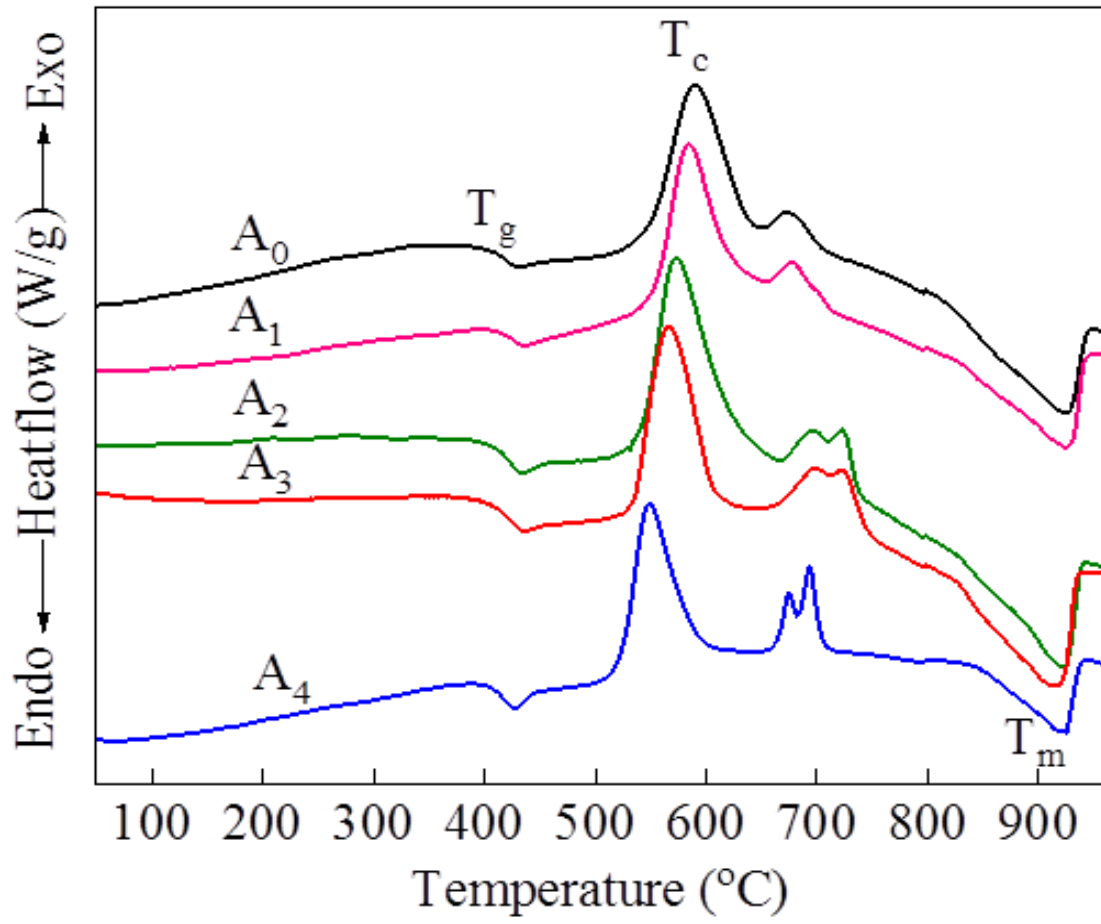


Fig. 4 DTA traces of heat treated  $\text{Li}_2\text{O-Pb}_3\text{O}_4\text{-SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

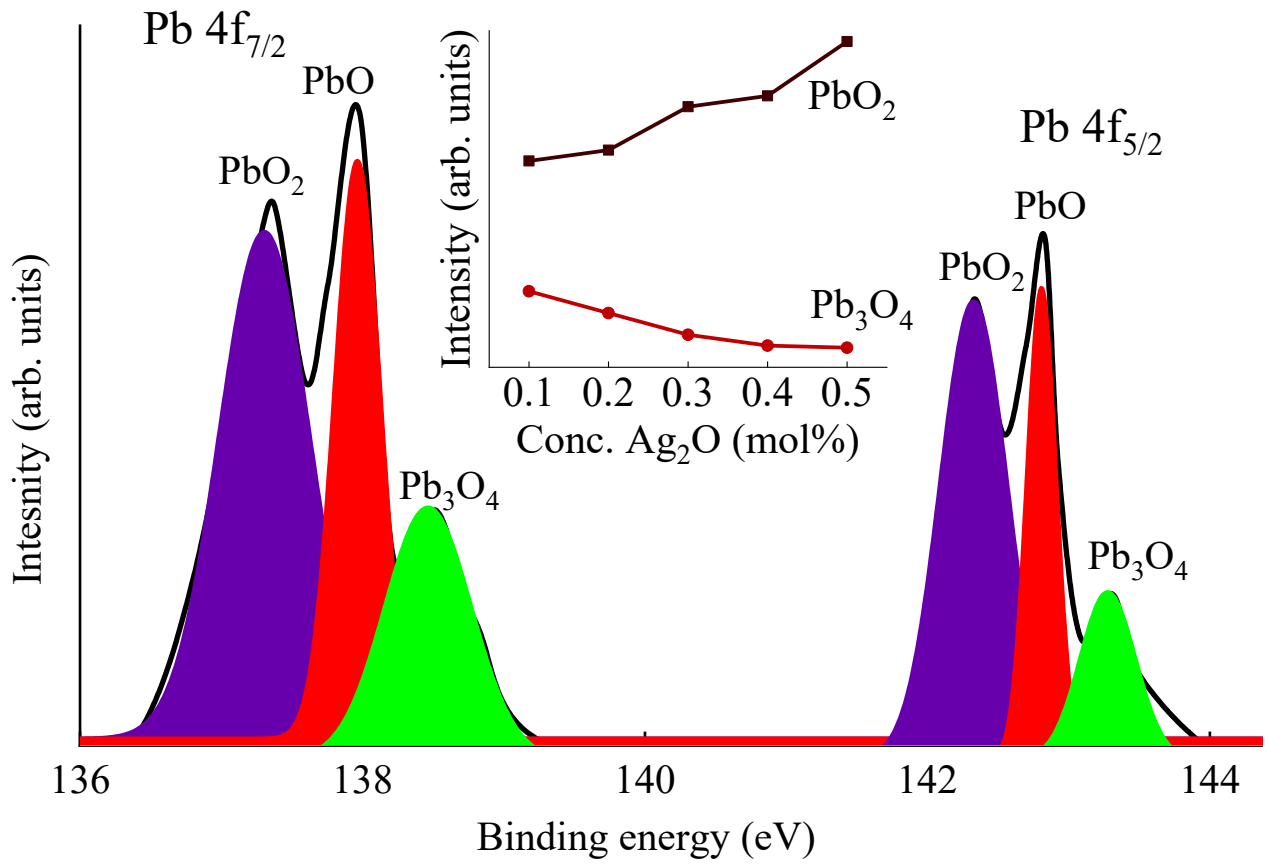


Fig. 5 Deconvoluted and fitted XPS spectra of PbO, PbO<sub>2</sub> and Pb<sub>3</sub>O<sub>4</sub> in the glass sample A<sub>2</sub> and inset represents the intensity variation of PbO<sub>2</sub> and Pb<sub>3</sub>O<sub>4</sub> (4f<sub>7/2</sub>) with the concentration of silver oxide.

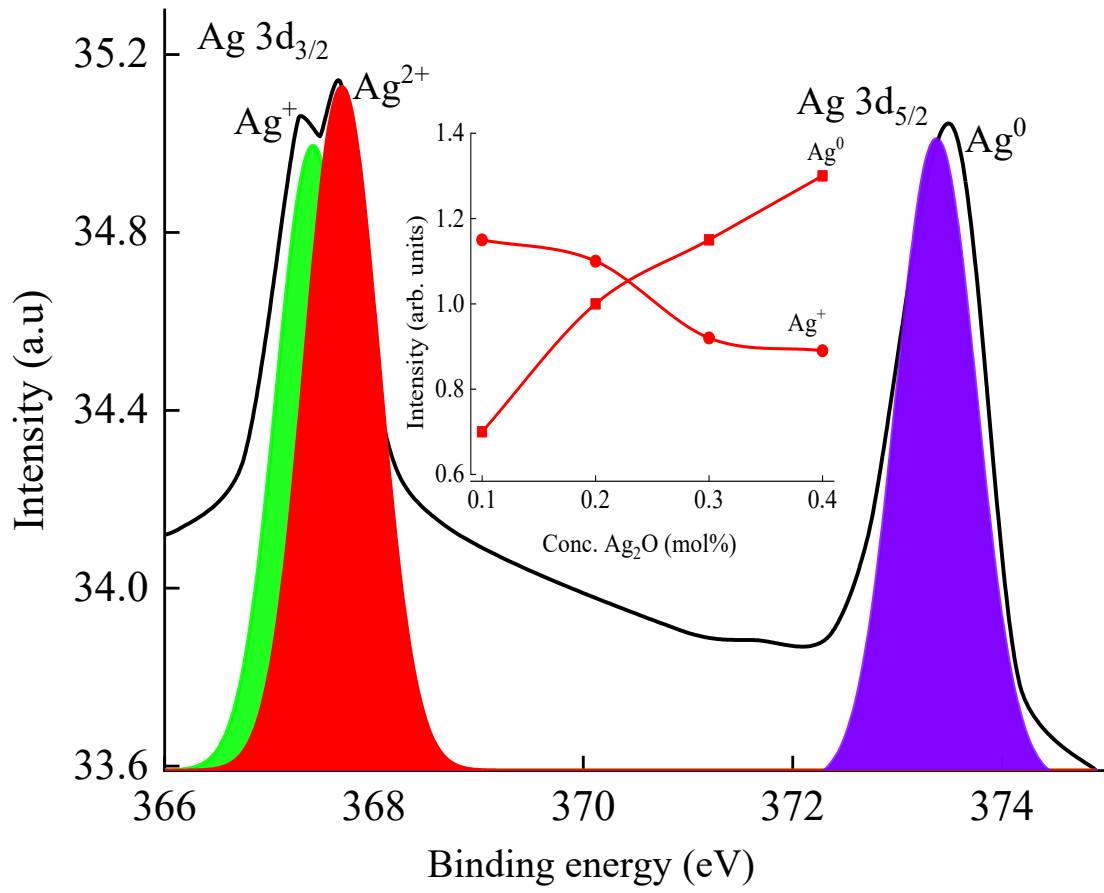


Fig. 6 Deconvoluted and fitted XPS spectra of Ag<sup>0</sup> and Ag<sup>+</sup> ions in the glass samples A<sub>2</sub> and the insets represents the intensity variation of Ag<sup>0</sup> and Ag<sup>+</sup> ions with the concentration of silver oxide.



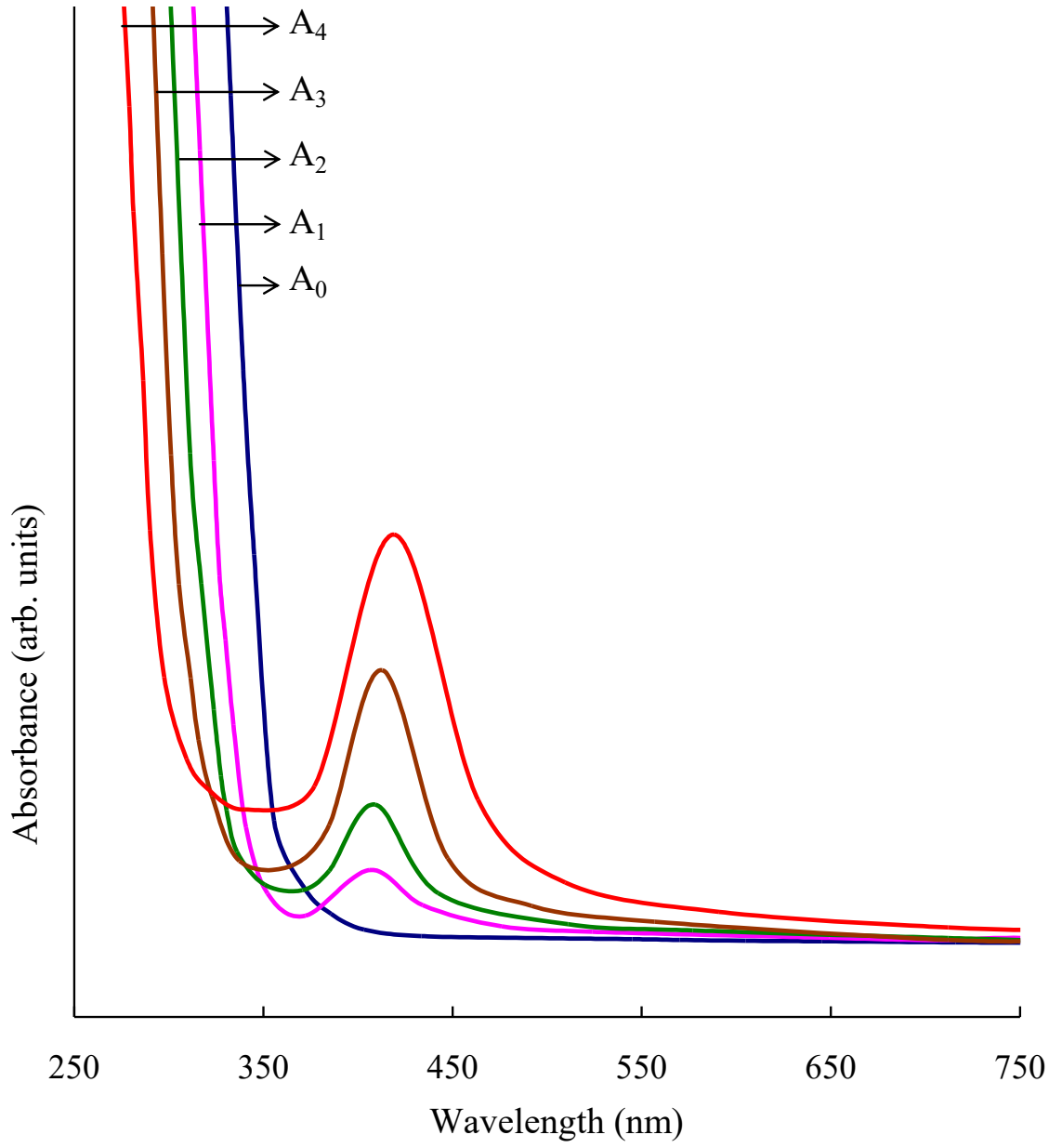


Fig. 7 Optical absorption spectra of  $\text{Li}_2\text{O-Pb}_3\text{O}_4\text{-SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$  recorded after heat treatment.

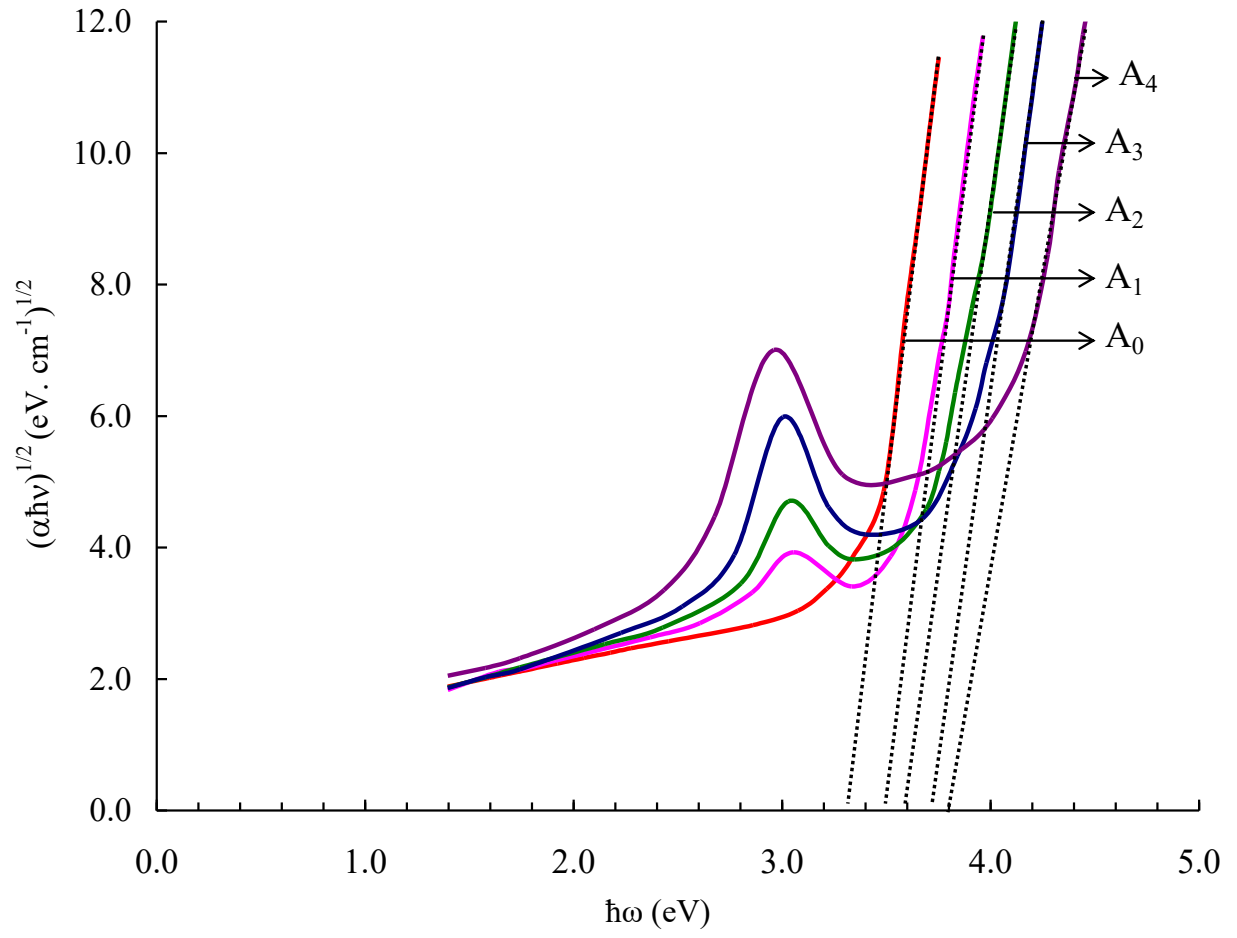


Fig. 8 Tauc plots of heat treated  $\text{Li}_2\text{O-Pb}_3\text{O}_4\text{-SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$ .

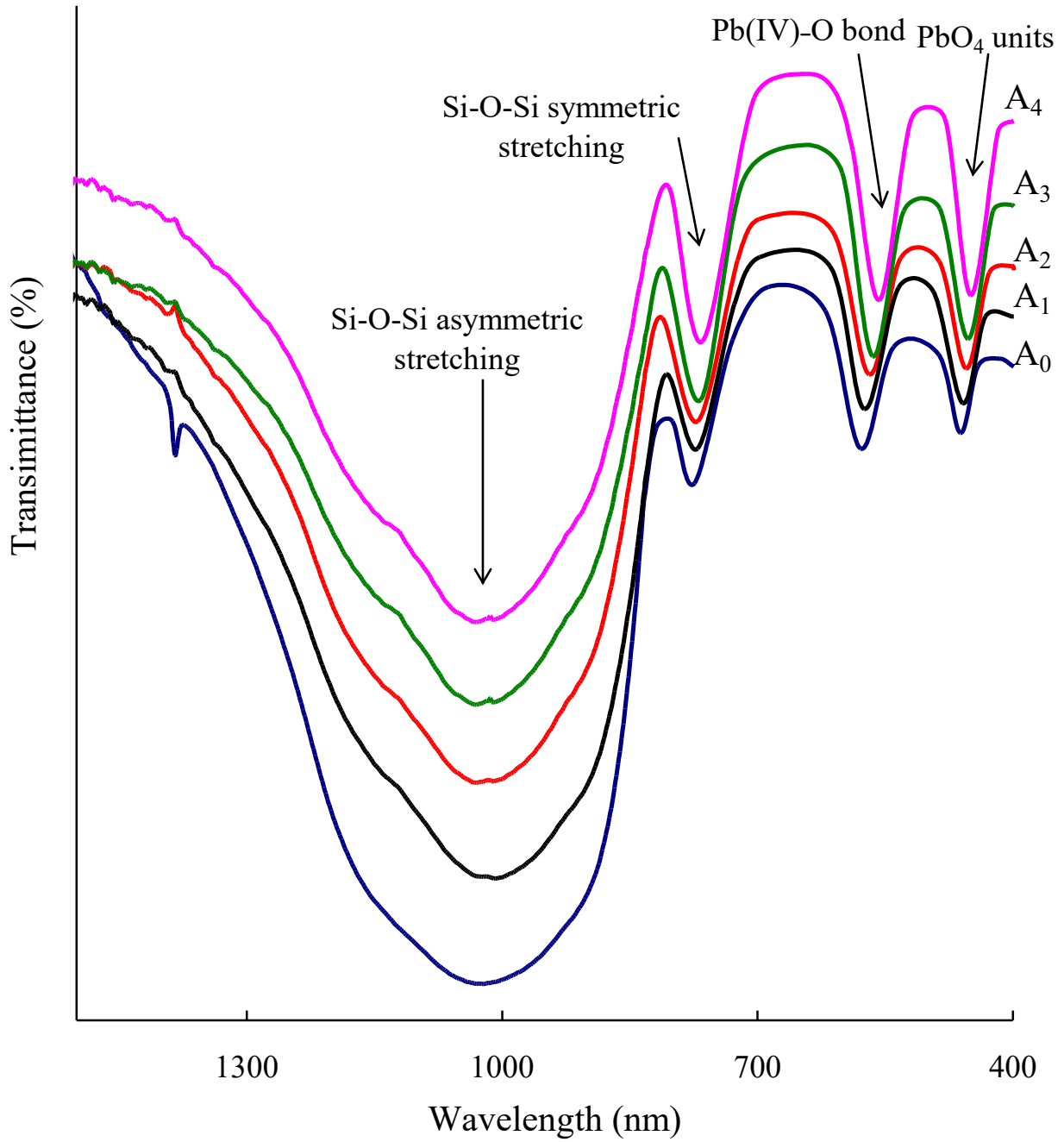


Fig. 9 FT-IR spectra of heat treated Li<sub>2</sub>O-Pb<sub>3</sub>O<sub>4</sub>-SiO<sub>2</sub> glasses doped with different concentrations of Ag<sub>2</sub>O.

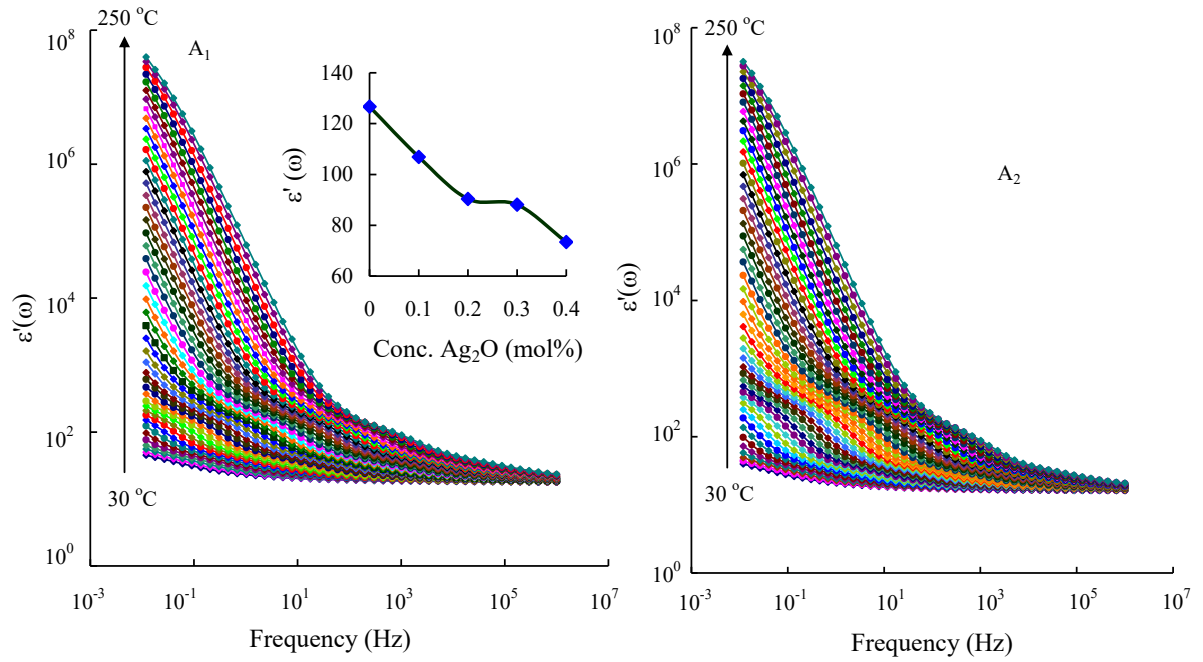


Fig. 10 (a) Variations of dielectric constant with frequency measured at different temperatures for the sample A<sub>1</sub> and A<sub>2</sub>.

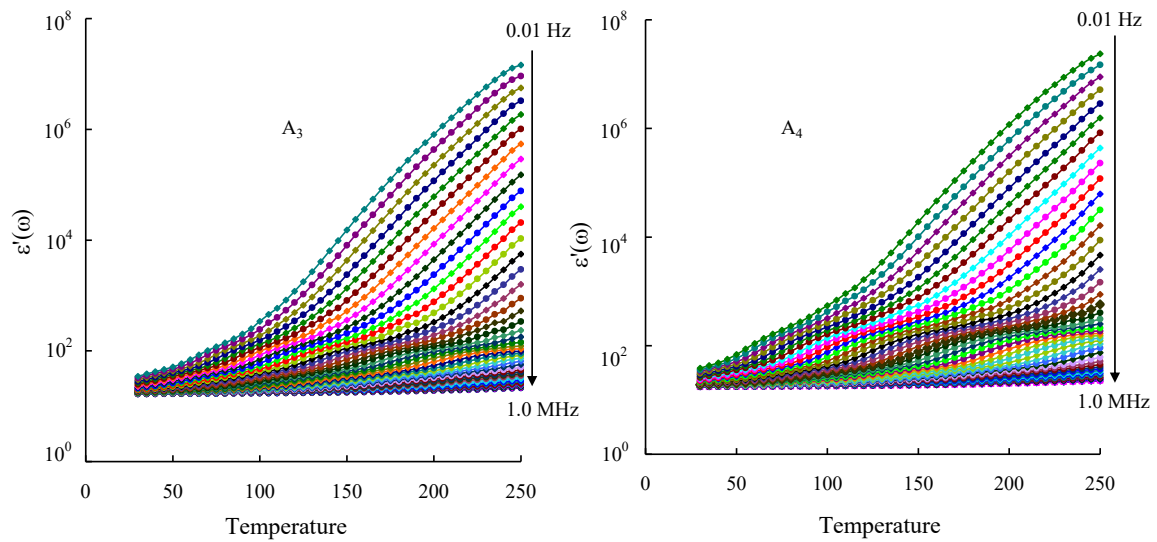


Fig. 10 (b) Variations of dielectric constant with temperature measured at different frequencies for the sample A<sub>3</sub> and A<sub>4</sub>.

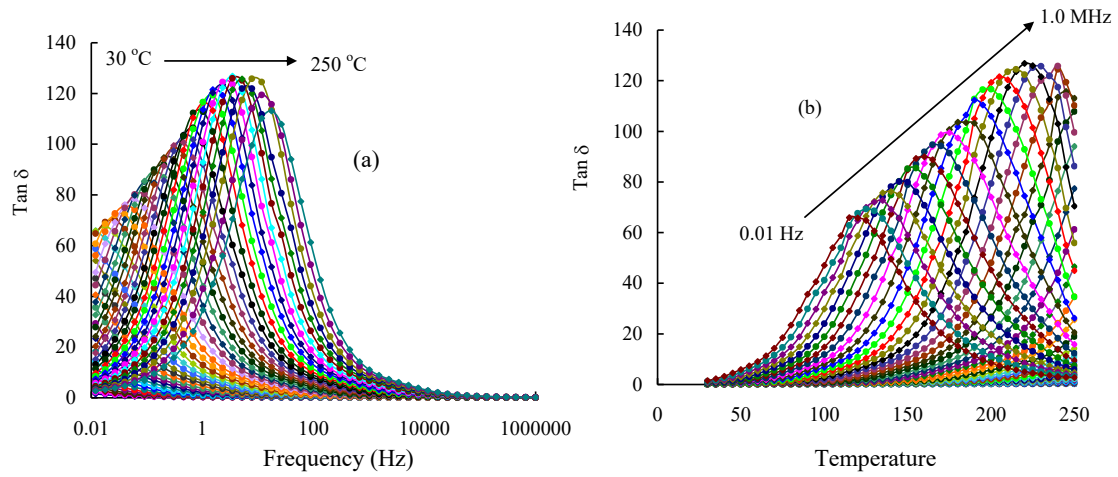


Fig. 11 Variations of dielectric loss (a) with frequency measured at different temperatures and (b) with temperatures for the sample A4.

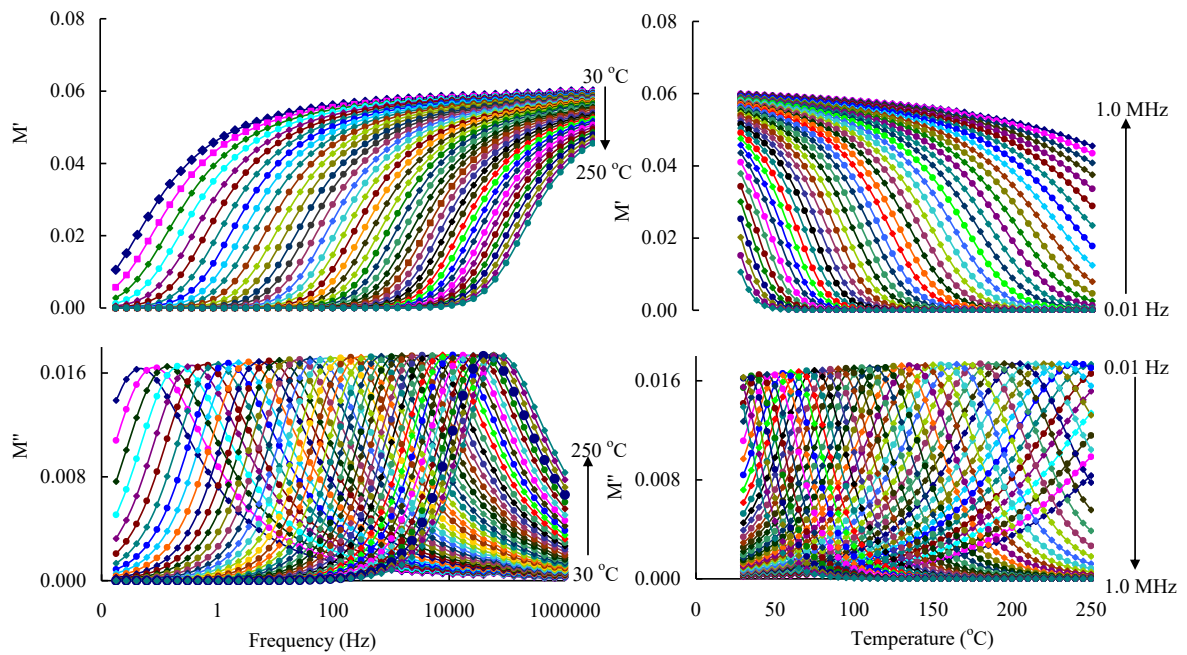


Fig. 12 Variations of electric moduli  $M'$  and  $M''$  (a) with frequency at different temperatures and (b) with temperature at different frequencies for the sample A3.

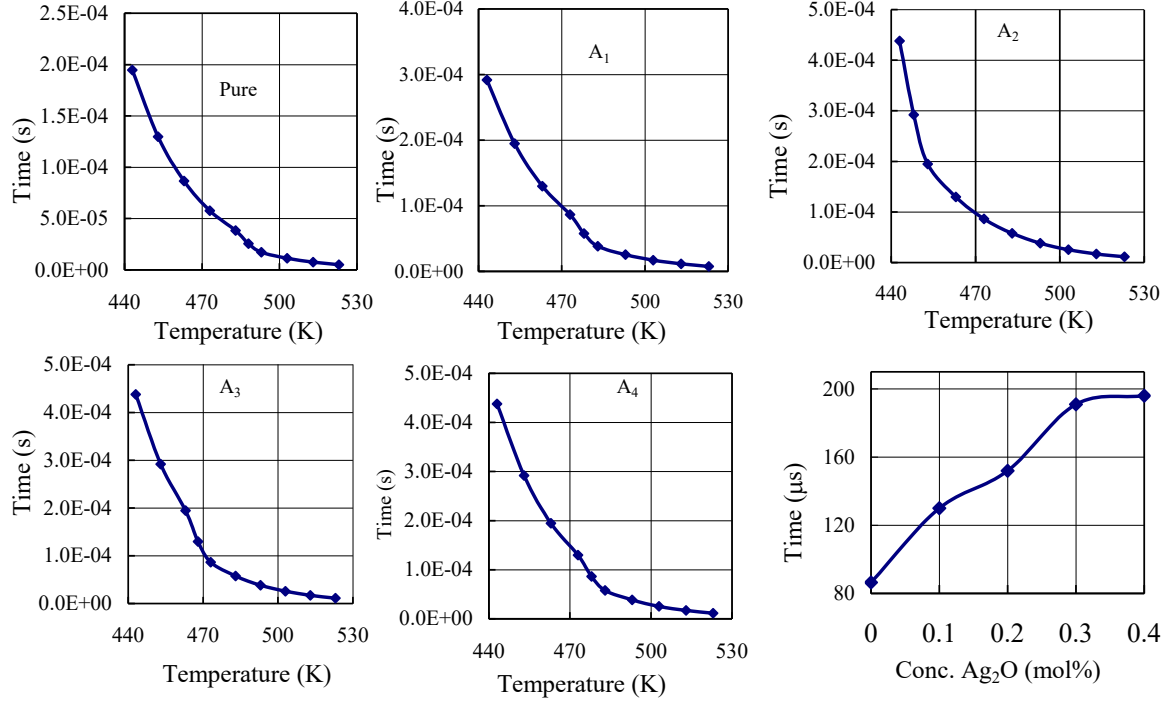


Fig. 13 The plots of relaxation times vs temperature for  $Li_2O-Pb_3O_4-SiO_2$  glasses doped with different concentrations of  $Ag_2O$ . In this figure the variations of  $\tau$  evaluated at 470 K with the concentration of  $Ag_2O$  is also presented.

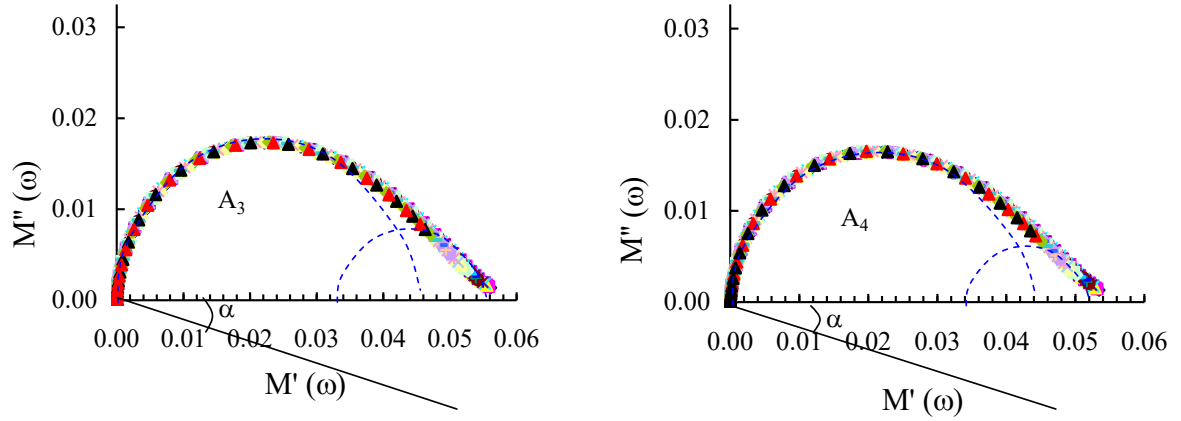


Fig. 14 Cole-cole plots for the samples  $A_3$  and  $A_4$ .

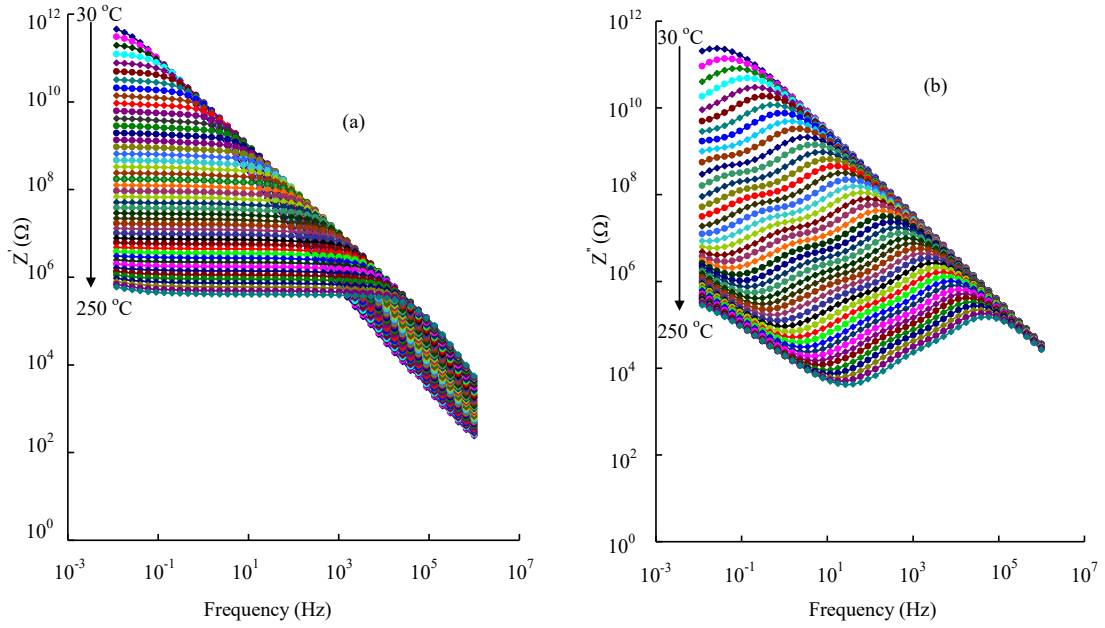


Fig. 15 Variations of  $Z'$  vs  $Z''$  with frequency at different temperatures for the samples  $A_2$ .

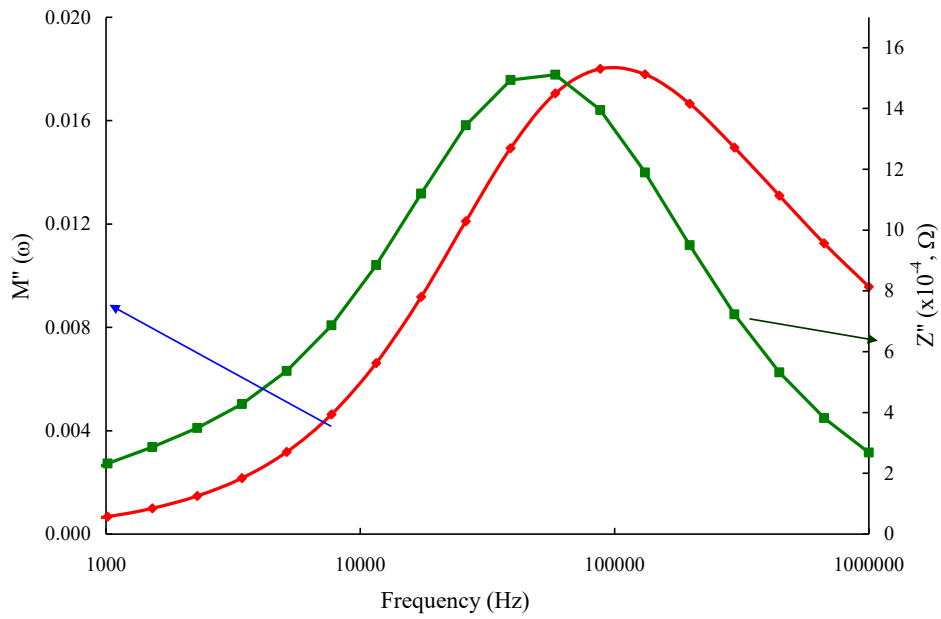


Fig. 15(c) Variation of  $M''(\omega)$  and  $Z''(\omega)$  with frequency at 200 °C for the glass  $A_2$ .

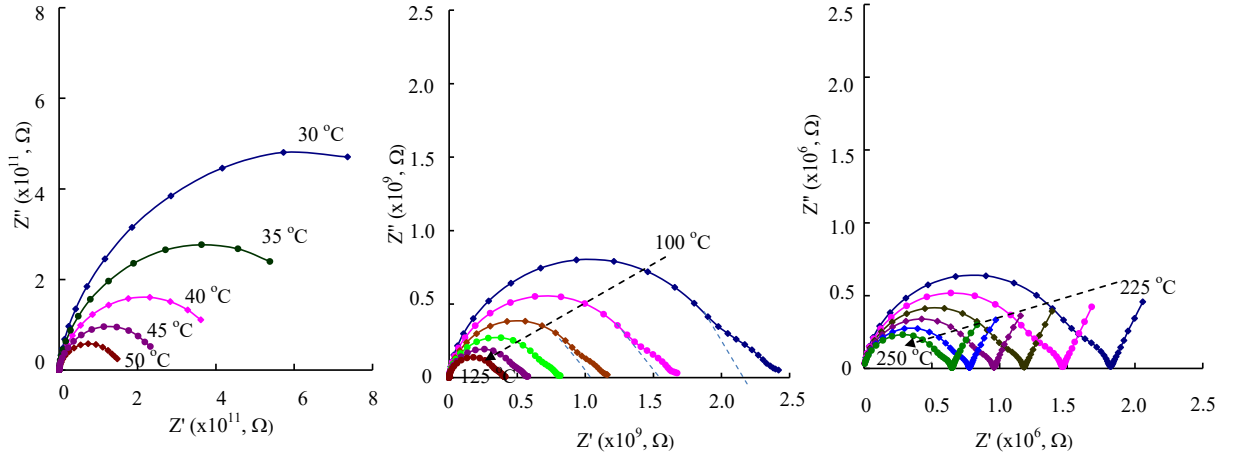


Fig. 16 (a) The plots of  $Z'$  vs  $Z''$  drawn in the temperature regions (b) 30-50 °C (b) 100-150 °C (c) 225-250 °C for the sample A<sub>4</sub>.

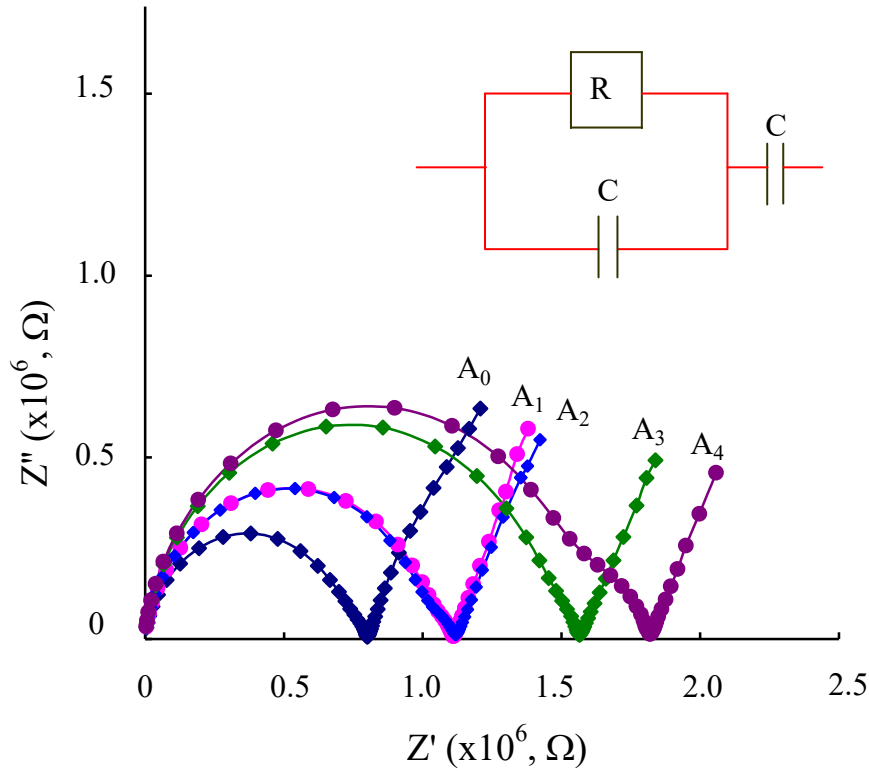


Fig. 16(d) The comparison plots of impedance of temperature of  $\text{Li}_2\text{O-Pb}_3\text{O}_4\text{-SiO}_2$  glasses doped with different concentrations of  $\text{Ag}_2\text{O}$  drawn at 250 °C. In the equivalent circuit diagram the parallel CR circuit accounts for semi-circular path of the diagram whereas additional series capacitor represents the inclined spur part.



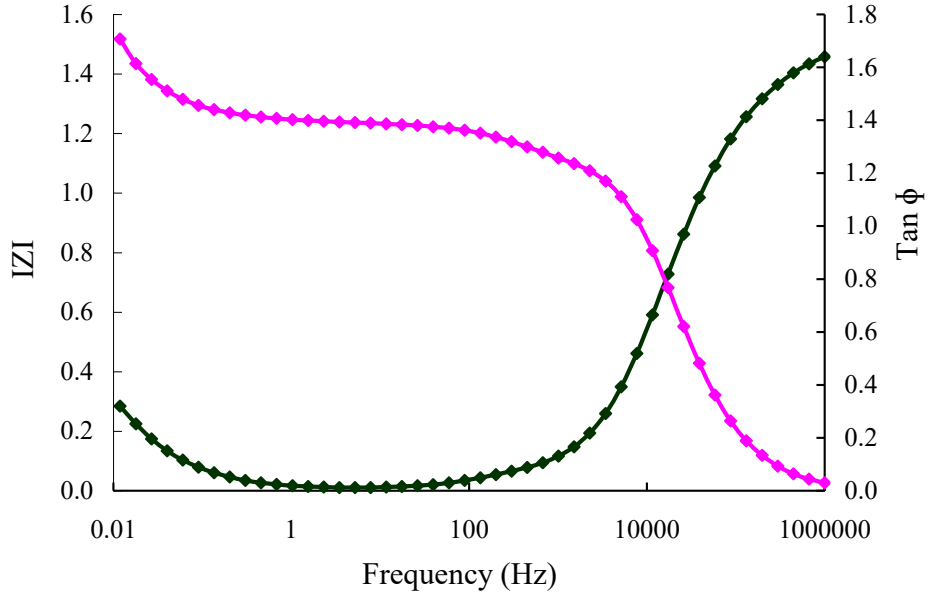


Fig.16(e) Bode plots (plots of total impedance  $|Z|$  and the phase angle,  $\tan \phi = Z''/Z'$  vs  $\omega$ ) for a selected glass sample A4 drawn at 200 °C.

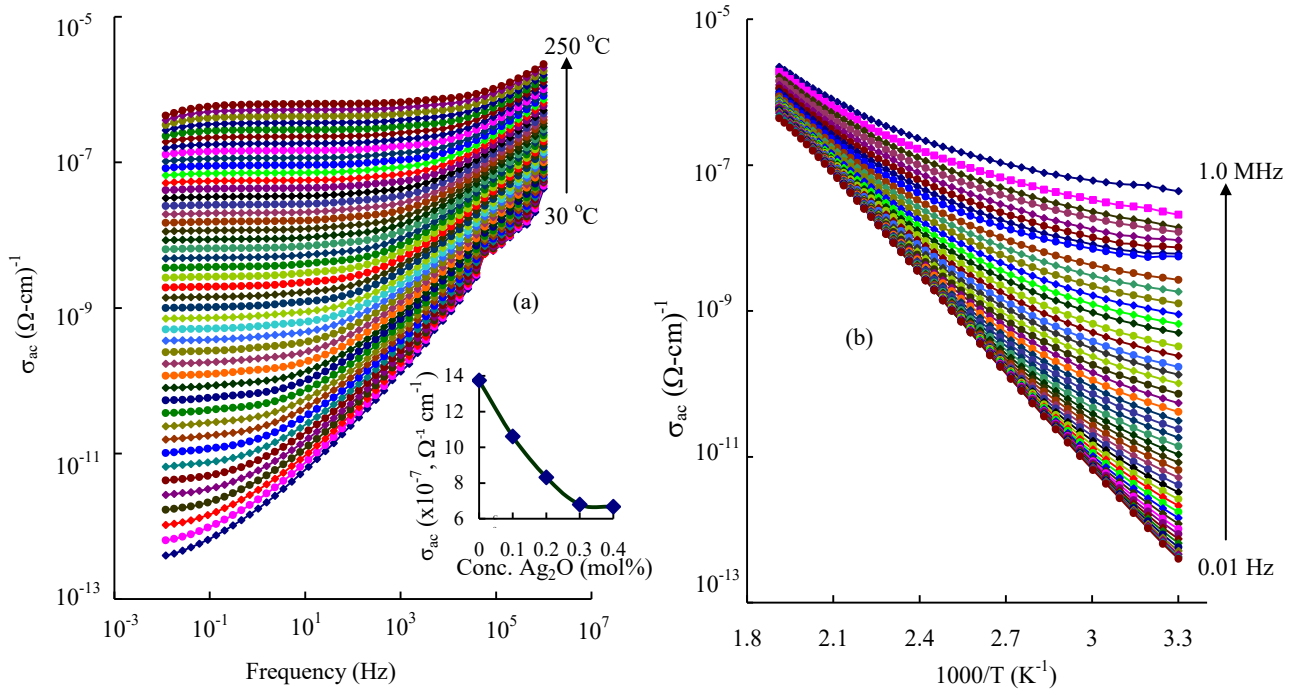


Fig. 17 The plots of  $\sigma_{ac}$  (a) vs frequency drawn at different temperatures and (b) vs  $1/T$  drawn at different frequencies for the glass samples A4. Inset represents the variation of  $\sigma_{ac}$  with concentration of Ag<sub>2</sub>O measured at 250°C and 1 kHz.

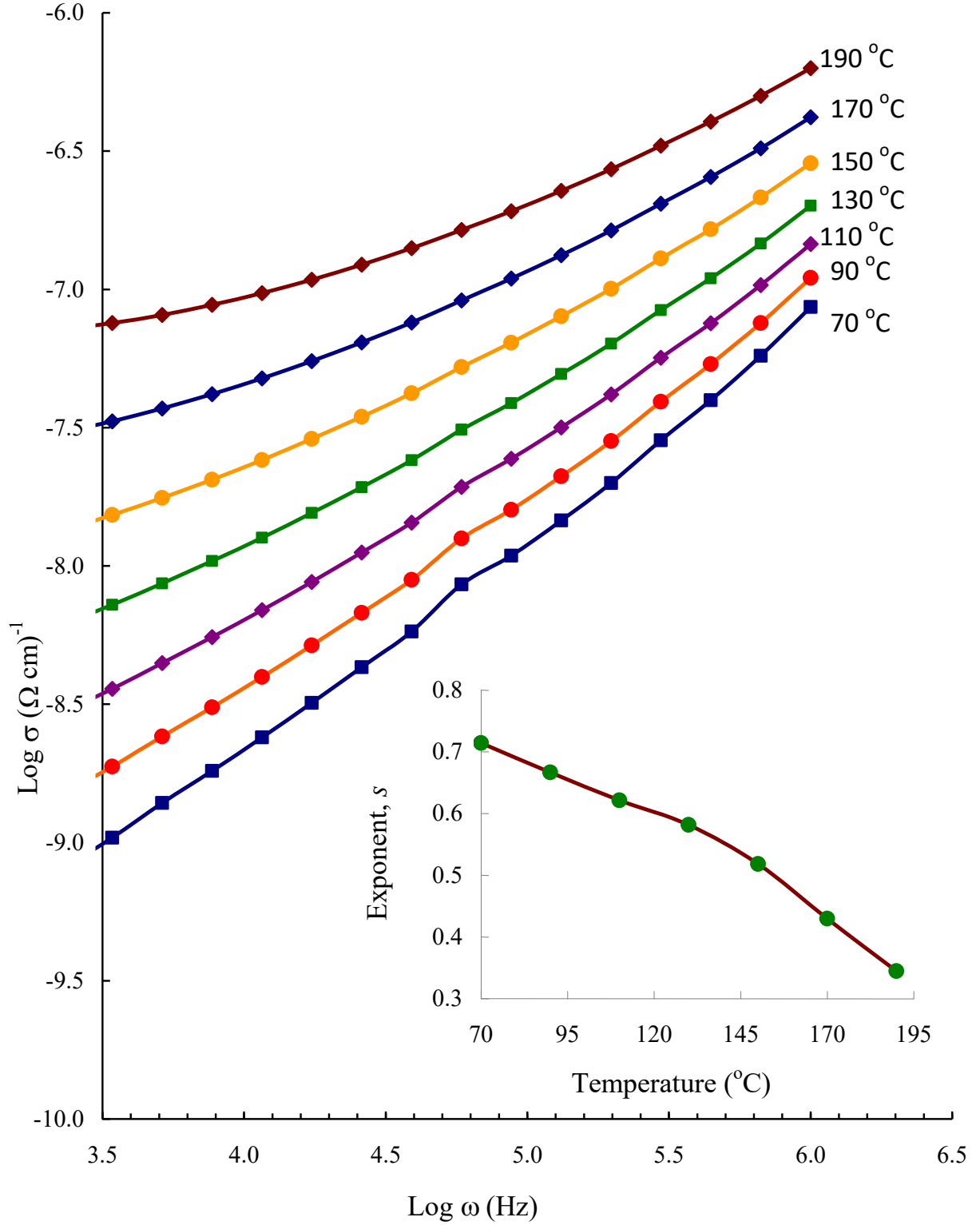


Fig. 18 Plots of  $\log \sigma_{ac}(\omega)$  versus  $\log \omega$  for the glass sample A<sub>2</sub> drawn at different temperatures and to evaluate frequency exponents.